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**MAE4700/5700**

**Finite Element Analysis for  
Mechanical and Aerospace Design**

**Cornell University, Fall 2009**

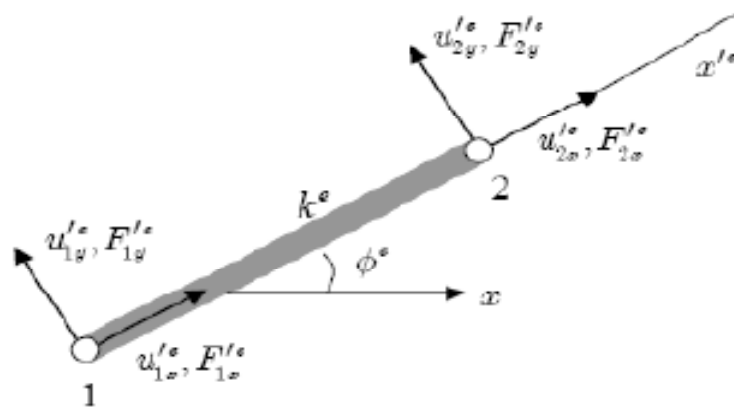
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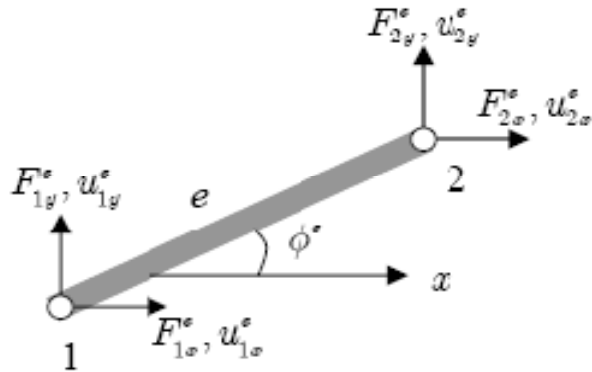
# Truss analysis

- Internal forces in a truss element act along the member  $F_{1y}^{(e)} = F_{2y}^{(e)} = 0$
- However, displacements at the nodes can have both components (x'- and y'-directions, in local coordinates).
- This is due to rotation



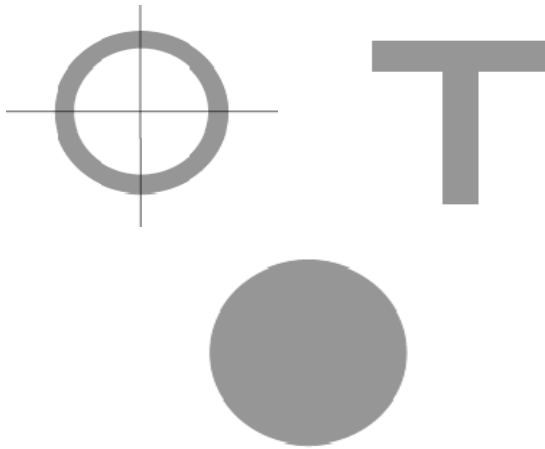
# Truss analysis

- To analyze a truss element in the **global coordinates  $x$  and  $y$** , you need to account for both components of displacement:

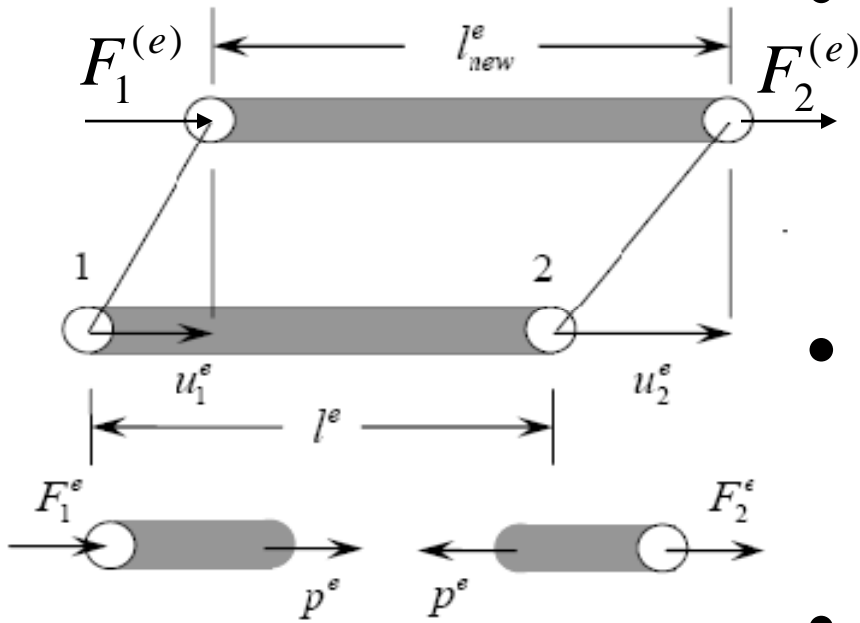


$$u_{1x}^{(e)}, u_{1y}^{(e)}, u_{2x}^{(e)}, u_{2y}^{(e)}$$

- Also note that the cross section of truss elements can vary as shown.



# Stiffness of a truss element



$$\begin{Bmatrix} F_1^{(e)} \\ F_2^{(e)} \end{Bmatrix} = \begin{bmatrix} k^{(e)} & -k^{(e)} \\ -k^{(e)} & k^{(e)} \end{bmatrix} \begin{Bmatrix} u_1^{(e)} \\ u_2^{(e)} \end{Bmatrix}$$

$$k^e = \frac{A^e E^e}{L^e}$$

- The internal force  $p^e$  in the truss is given (see free body diagram) as:

$$p^e = F_2^{(e)} = -F_1^{(e)} = A^e \sigma^e$$

- Assuming elastic deformations:

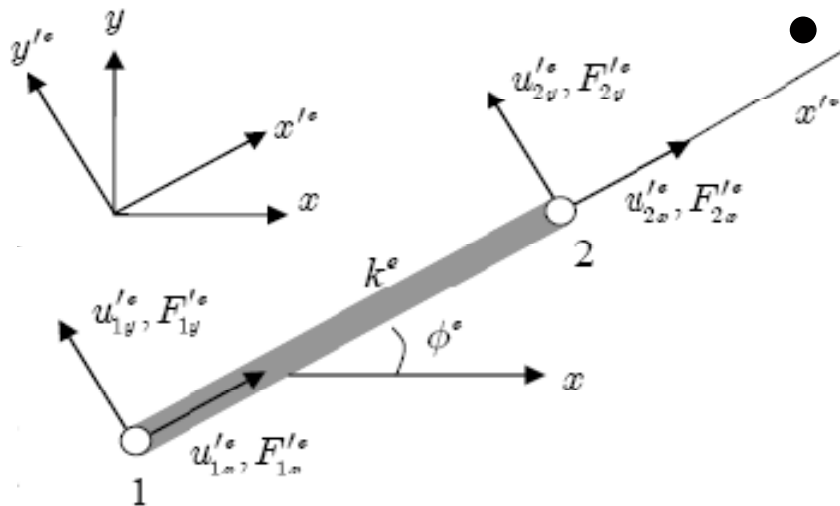
$$p^e = F_2^{(e)} = -F_1^{(e)} = A^e E^e \varepsilon^e$$

- The (small) strain is given as:

$$\varepsilon^e = \frac{u_2^e - u_1^e}{L^e}$$

- Finally:  $F_2^{(e)} = -F_1^{(e)} = \frac{A^e E^e}{L^e} (u_2^e - u_1^e) = k^e (u_2^e - u_1^e)$

# Truss element stiffness in local coordinates



- We can re-write the element stiffness equations as:

$$\begin{Bmatrix} F_{1x}^{(e)} \\ F_{1y}^{(e)} \\ F_{2x}^{(e)} \\ F_{2y}^{(e)} \end{Bmatrix} = \begin{bmatrix} k^{(e)} & -k^{(e)} \\ -k^{(e)} & k^{(e)} \end{bmatrix} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{2x}^{(e)} \end{Bmatrix}$$

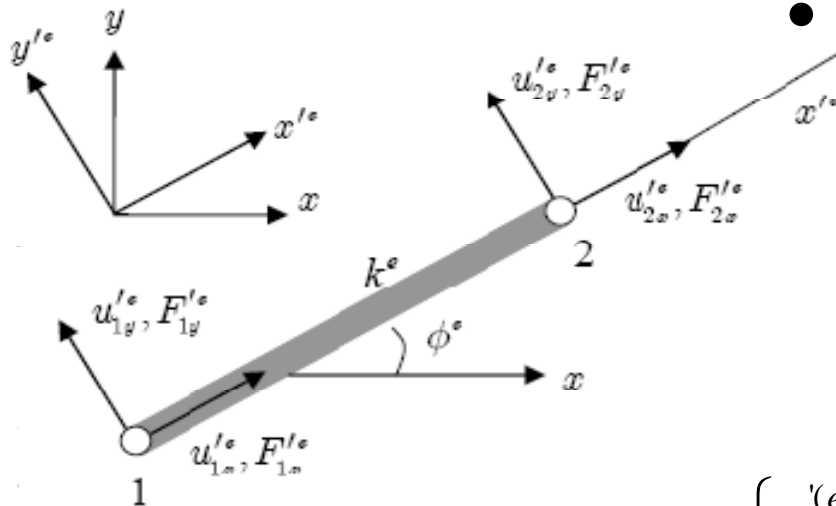
$$\begin{Bmatrix} F_{1x}^{(e)} \\ F_{1y}^{(e)} \\ F_{2x}^{(e)} \\ F_{2y}^{(e)} \end{Bmatrix} = k^{(e)} \underbrace{\begin{bmatrix} 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}}_{[K^{(e)}]} \underbrace{\begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix}}_{\{d^{(e)}\}}$$

Notice (as it should be) that:

$$F_{1y}^{(e)} = F_{2y}^{(e)} = 0$$

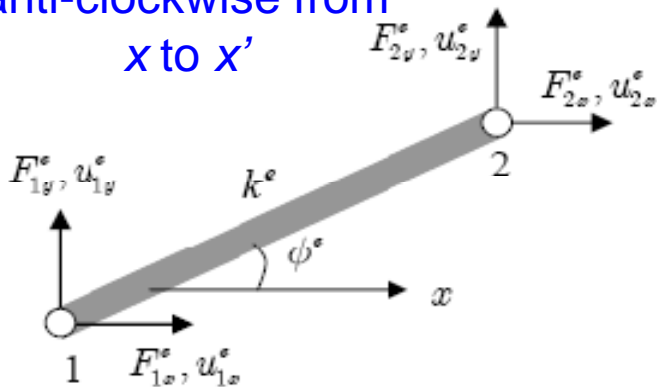
$$\{F^{(e)}\} = [K^{(e)}] \{d^{(e)}\}$$

# Element stiffness in global coordinates



- We need to be able to transform displacements from the  $x'$  and  $y'$  axes to displacements along the  $x$  and  $y$  axes. We start with the reverse:

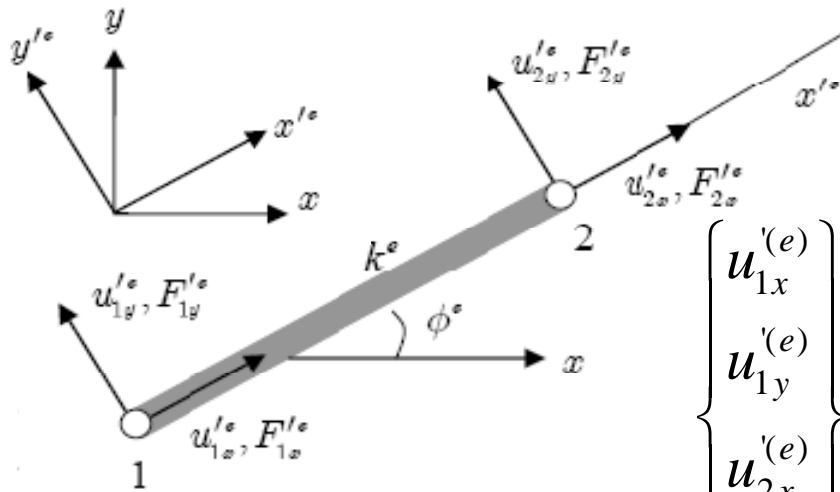
The angle  $\phi^e$  is measured anti-clockwise from  $x$  to  $x'$



$$\begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix} = \begin{bmatrix} \cos \phi^e & \sin \phi^e & 0 & 0 \\ -\sin \phi^e & \cos \phi^e & 0 & 0 \\ 0 & 0 & \cos \phi^e & \sin \phi^e \\ 0 & 0 & -\sin \phi^e & \cos \phi^e \end{bmatrix} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix}$$

Transformation matrix  $T^{(e)}$

# Coordinate transformation



$$\{d^{'e}\} = [T^e] \{d^e\}$$

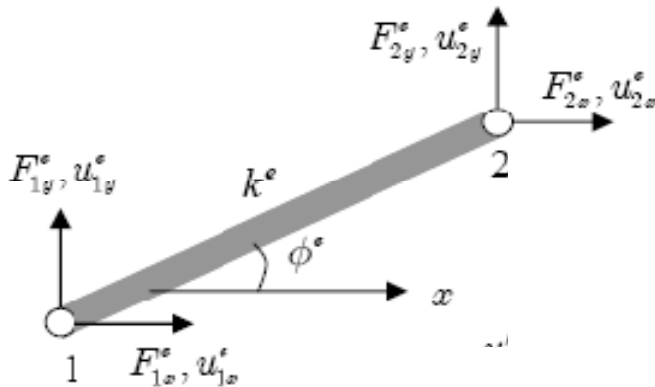
$$\begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix} = \begin{bmatrix} \cos \phi^e & \sin \phi^e & 0 & 0 \\ -\sin \phi^e & \cos \phi^e & 0 & 0 \\ 0 & 0 & \cos \phi^e & \sin \phi^e \\ 0 & 0 & -\sin \phi^e & \cos \phi^e \end{bmatrix} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix}$$

$\underbrace{\hspace{10em}}_{[T^{(e)}]}$

Note that :

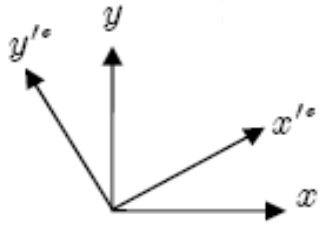
$$[T^e]^T [T^e] = [I]$$

$$\{d^e\} = [T^e]^T \{d^{'e}\}$$





# Coordinate transformation

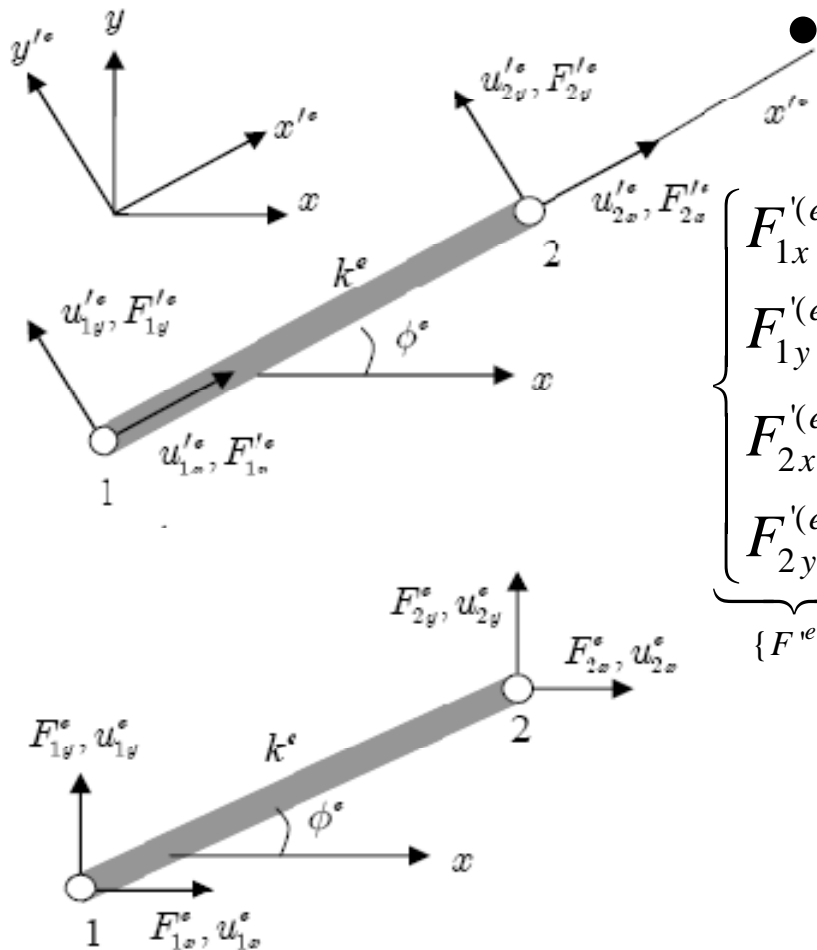


Verify that:  $[T^e]^T [T^e] = [I]_{4 \times 4}$

$$\underbrace{\begin{bmatrix} \cos \phi^e & -\sin \phi^e & 0 & 0 \\ \sin \phi^e & \cos \phi^e & 0 & 0 \\ 0 & 0 & \cos \phi^e & -\sin \phi^e \\ 0 & 0 & \sin \phi^e & \cos \phi^e \end{bmatrix}}_{T^{eT}} \underbrace{\begin{bmatrix} \cos \phi^e & \sin \phi^e & 0 & 0 \\ -\sin \phi^e & \cos \phi^e & 0 & 0 \\ 0 & 0 & \cos \phi^e & \sin \phi^e \\ 0 & 0 & -\sin \phi^e & \cos \phi^e \end{bmatrix}}_{T^e} =$$

$$\underbrace{\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}}_{I_{4 \times 4}}$$

# Stiffness of a truss element



• Similarly for the forces:

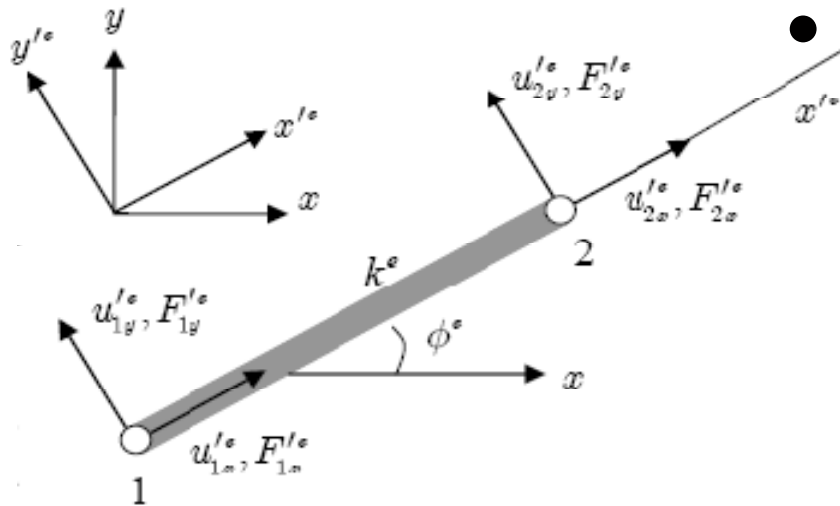
$$\{F'^e\} = [T^e] \{F^e\}$$

$$\begin{Bmatrix} F'_{1x} \\ F'_{1y} \\ F'_{2x} \\ F'_{2y} \end{Bmatrix} = \begin{bmatrix} \cos \phi^e & \sin \phi^e & 0 & 0 \\ -\sin \phi^e & \cos \phi^e & 0 & 0 \\ 0 & 0 & \cos \phi^e & \sin \phi^e \\ 0 & 0 & -\sin \phi^e & \cos \phi^e \end{bmatrix} \begin{Bmatrix} F_{1x} \\ F_{1y} \\ F_{2x} \\ F_{2y} \end{Bmatrix}$$

$\underbrace{\hspace{15em}}_{[T^{(e)}]}$

$$\{F^e\} = [T^e]^T \{F'^e\}$$

# Stiffness of a truss element



- Using  $\{F^{(e)}\} = [K^{(e)}] \{d^{(e)}\}$  and the transformation equations,  $\{F^{(e)}\} = [T^{(e)}] \{F^{(e)}\}$   
 $\{d^{(e)}\} = [T^{(e)}] \{d^{(e)}\}$

we can write the stiffness in the  $x, y$  system as follows:

$$\{F^{(e)}\} = [K^{(e)}] \{d^{(e)}\} \Rightarrow$$

$$[T^{(e)}] \{F^{(e)}\} = [K^{(e)}] [T^{(e)}] \{d^{(e)}\} \Rightarrow \{F^{(e)}\} = \underbrace{[T^{(e)}]^T [K^{(e)}] [T^{(e)}]}_{[K^{(e)}]} \{d^{(e)}\} \Rightarrow$$

$$[K^{(e)}] = [T^{(e)}]^T [K^{(e)}] [T^{(e)}]$$

# Truss element stiffness

$$[K^{(e)}] = [T^{(e)}]^T [K'^{(e)}] [T^{(e)}]$$

$$[K'^{(e)}] = k^{(e)} \begin{bmatrix} 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}, \quad [T^{(e)}] = \begin{bmatrix} \cos \phi^e & \sin \phi^e & 0 & 0 \\ -\sin \phi^e & \cos \phi^e & 0 & 0 \\ 0 & 0 & \cos \phi^e & \sin \phi^e \\ 0 & 0 & -\sin \phi^e & \cos \phi^e \end{bmatrix}$$

$$[K^{(e)}] = k^{(e)} \begin{bmatrix} \cos^2 \phi^e & \sin \phi^e \cos \phi^e & -\cos^2 \phi^e & -\sin \phi^e \cos \phi^e \\ \sin \phi^e \cos \phi^e & \sin^2 \phi^e & -\sin \phi^e \cos \phi^e & -\sin^2 \phi^e \\ -\cos^2 \phi^e & -\sin \phi^e \cos \phi^e & \cos^2 \phi^e & \sin \phi^e \cos \phi^e \\ -\sin \phi^e \cos \phi^e & -\sin^2 \phi^e & \sin \phi^e \cos \phi^e & \sin^2 \phi^e \end{bmatrix}$$

# Truss element stiffness

$$\begin{Bmatrix} F_{1x}^{(e)} \\ F_{1y}^{(e)} \\ F_{2x}^{(e)} \\ F_{2y}^{(e)} \end{Bmatrix} = k^{(e)} \begin{bmatrix} \cos^2 \phi^e & \sin \phi^e \cos \phi^e & -\cos^2 \phi^e & -\sin \phi^e \cos \phi^e \\ \sin \phi^e \cos \phi^e & \sin^2 \phi^e & -\sin \phi^e \cos \phi^e & -\sin^2 \phi^e \\ \hline -\cos^2 \phi^e & -\sin \phi^e \cos \phi^e & \cos^2 \phi^e & \sin \phi^e \cos \phi^e \\ -\sin \phi^e \cos \phi^e & -\sin^2 \phi^e & \sin \phi^e \cos \phi^e & \sin^2 \phi^e \end{bmatrix} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix}$$

- Note the 2x2 symmetric submatrix structure
- This implies that you can reverse the numbering of nodes (1 and 2) without any changes in the element stiffness.

$$\begin{Bmatrix} F_{2x}^{(e)} \\ F_{2y}^{(e)} \\ F_{1x}^{(e)} \\ F_{1y}^{(e)} \end{Bmatrix} = k^{(e)} \begin{bmatrix} \cos^2 \phi^e & \sin \phi^e \cos \phi^e & -\cos^2 \phi^e & -\sin \phi^e \cos \phi^e \\ \sin \phi^e \cos \phi^e & \sin^2 \phi^e & -\sin \phi^e \cos \phi^e & -\sin^2 \phi^e \\ \hline -\cos^2 \phi^e & -\sin \phi^e \cos \phi^e & \cos^2 \phi^e & \sin \phi^e \cos \phi^e \\ -\sin \phi^e \cos \phi^e & -\sin^2 \phi^e & \sin \phi^e \cos \phi^e & \sin^2 \phi^e \end{bmatrix} \begin{Bmatrix} u_{2x}^{(e)} \\ u_{2y}^{(e)} \\ u_{1x}^{(e)} \\ u_{1y}^{(e)} \end{Bmatrix}$$

# Assembly process

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- The assembly process is identical to the one discussed for 'spring structures' and it will not be repeated here in its general form (no need to show at this point complicated looking matrix operations).
- We will however provide soon a simple example demonstrating this assembly process.

# Generalizing the application of essential BCs

- You already have seen through an example how essential boundary conditions are applied to the global system of eqs:  $[K]\{d\} = \{F\}$
- In essence, we partition the stiffness matrix in a way that separates known from unknown degrees of freedom as follows:

$$\begin{bmatrix} K_E & K_{EF} \\ K_{EF}^T & K_F \end{bmatrix} \begin{Bmatrix} \bar{d}_E \\ d_F \end{Bmatrix} = \begin{Bmatrix} \bar{f}_E \\ f_F \end{Bmatrix}$$

$\bar{d}_E$  : Known displacements

$d_F$  : Unknown displacements

$f_F$  : Applied (known) forces

$\bar{f}_E$  : Unknown reaction forces

corresponding to  
nodes/directions

with prescribed displacement

# Generalizing the application of essential BCs

$$\begin{bmatrix} K_E & K_{EF} \\ K_{EF}^T & K_F \end{bmatrix} \begin{Bmatrix} \bar{d}_E \\ d_F \end{Bmatrix} = \begin{Bmatrix} \bar{f}_E \\ f_F \end{Bmatrix} \Rightarrow \begin{aligned} K_E \bar{d}_E + K_{EF} d_F &= \bar{f}_E \\ K_{EF}^T \bar{d}_E + K_F d_F &= f_F \end{aligned}$$

- The unknown displacements are obtained from the 2<sup>nd</sup> eq. as:

$$K_{EF}^T \bar{d}_E + K_F d_F = f_F \Rightarrow d_F = K_F^{-1} (f_F - K_{EF}^T \bar{d}_E)$$

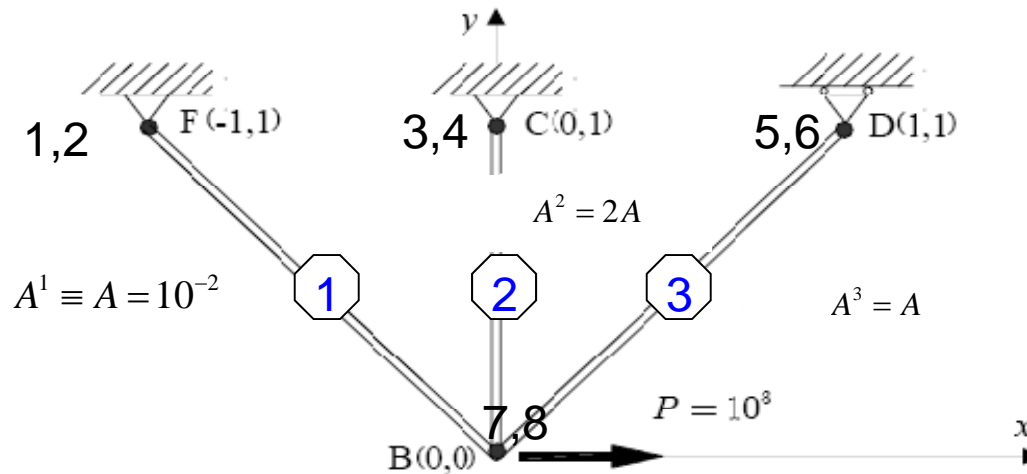
- With known  $d_F$ , we can return to the 1<sup>st</sup> eq. to compute the reaction forces:

$$\bar{f}_E = K_E \bar{d}_E + K_{EF} d_F$$

- Note that the matrix  $K_F$  is **symmetric and positive definite**, so a solution for  $d_F$  always exists!



# A truss example

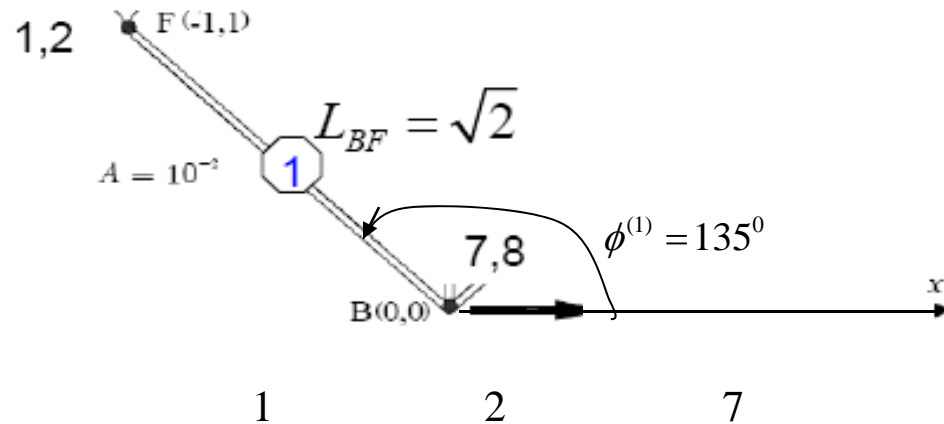


Note that point D is free to move in the x direction

$$E = 10^7 \text{ Pa}$$

- Construct the global stiffness matrix and load vector
- Partition the matrices and solve for the unknown displacements at point B, and displacement in x direction at point D.
- Find the stresses in the three bars
- Find the reactions at C, D and F

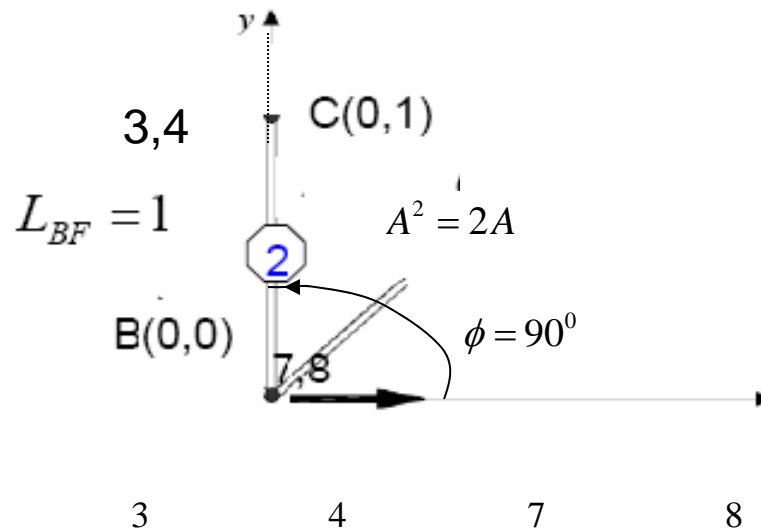
# A truss example: Element 1



$$[K^{(1)}] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 1/2 & -1/2 & -1/2 & 1/2 \\ -1/2 & 1/2 & 1/2 & -1/2 \\ -1/2 & 1/2 & 1/2 & -1/2 \\ 1/2 & -1/2 & -1/2 & 1/2 \end{bmatrix} \begin{matrix} 1 \\ 2 \\ 7 \\ 8 \end{matrix}$$

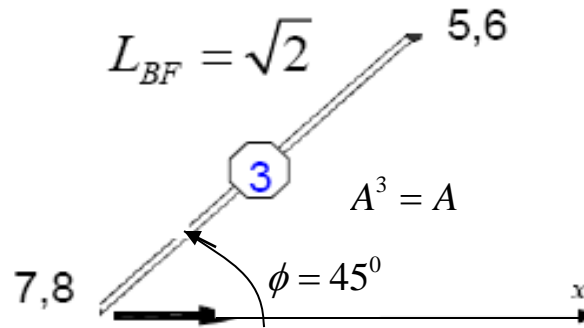
Note: Recall that you can number the corresponding global nodes in the sequence 7 8 1 2 without any changes in  $[K^{(1)}]$ .

# A truss example: Element 2



$$[K^{(2)}] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 2\sqrt{2} & 0 & -2\sqrt{2} \\ 0 & 0 & 0 & 0 \\ 0 & -2\sqrt{2} & 0 & 2\sqrt{2} \end{bmatrix} \begin{matrix} 3 \\ 4 \\ 7 \\ 8 \end{matrix}$$

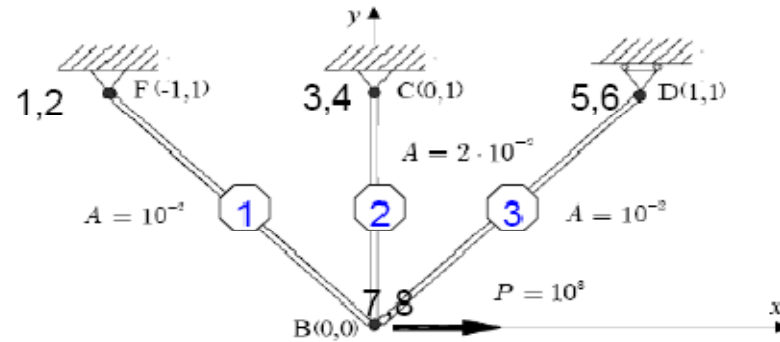
# A truss example: Element 3



$$[K^{(3)}] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 1/2 & 1/2 & -1/2 & -1/2 \\ 1/2 & 1/2 & -1/2 & -1/2 \\ -1/2 & -1/2 & 1/2 & 1/2 \\ -1/2 & -1/2 & 1/2 & 1/2 \end{bmatrix} \begin{matrix} 5 \\ 6 \\ 7 \\ 8 \end{matrix}$$

# A truss example: Assembly (element 1)

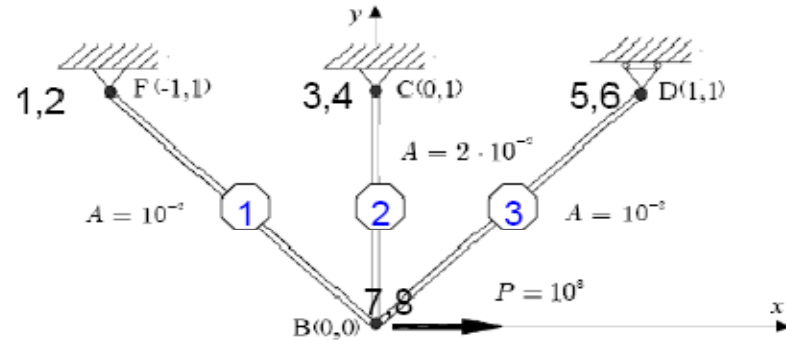
$$[K^{(1)}] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 1/2 & -1/2 & -1/2 & 1/2 \\ -1/2 & 1/2 & 1/2 & -1/2 \\ -1/2 & 1/2 & 1/2 & -1/2 \\ 1/2 & -1/2 & -1/2 & 1/2 \end{bmatrix} \begin{matrix} 1 \\ 2 \\ 7 \\ 8 \end{matrix}$$



$$[K] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 1/2 & -1/2 & 0 & 0 & 0 & 0 & -1/2 & -1/2 \\ -1/2 & 1/2 & 0 & 0 & 0 & 0 & 1/2 & -1/2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -1/2 & 1/2 & 0 & 0 & 0 & 0 & 1/2 & -1/2 \\ 1/2 & -1/2 & 0 & 0 & 0 & 0 & -1/2 & 1/2 \end{bmatrix} \begin{matrix} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \\ 8 \end{matrix}$$

# A truss example: Assembly (element 2)

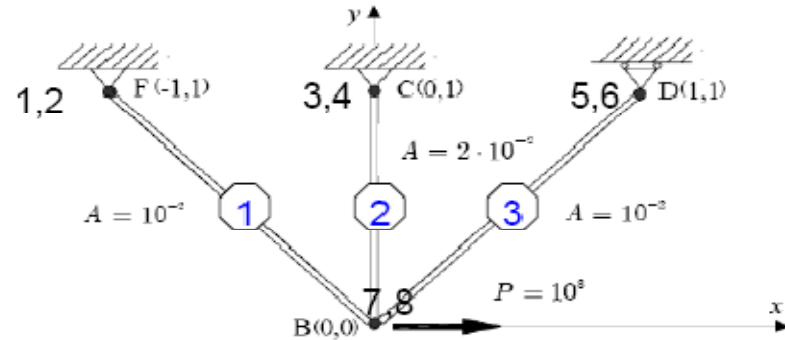
$$[K^{(2)}] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 2\sqrt{2} & 0 & -2\sqrt{2} \\ 0 & 0 & 0 & 0 \\ 0 & -2\sqrt{2} & 0 & 2\sqrt{2} \end{bmatrix} \begin{matrix} 3 \\ 4 \\ 7 \\ 8 \end{matrix}$$



$$[K] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 1/2 & -1/2 & 0 & 0 & 0 & 0 & -1/2 & -1/2 \\ -1/2 & 1/2 & 0 & 0 & 0 & 0 & 1/2 & -1/2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 2\sqrt{2} & 0 & 0 & 0 & -2\sqrt{2} \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -1/2 & 1/2 & 0 & 0 & 0 & 0 & 1/2+0 & -1/2+0 \\ 1/2 & -1/2 & 0 & -2\sqrt{2} & 0 & 0 & -1/2 & 1/2+2\sqrt{2} \end{bmatrix} \begin{matrix} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \\ 8 \end{matrix}$$

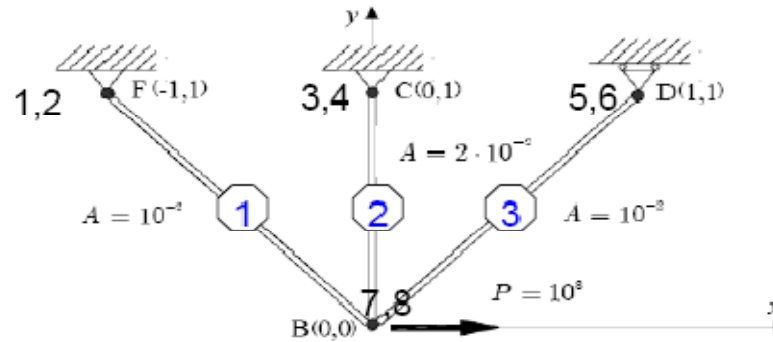
# A truss example: Assembly (element 3)

$$[K^{(3)}] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 1/2 & 1/2 & -1/2 & -1/2 \\ 1/2 & 1/2 & -1/2 & -1/2 \\ -1/2 & -1/2 & 1/2 & 1/2 \\ -1/2 & -1/2 & 1/2 & 1/2 \end{bmatrix} \begin{matrix} 5 \\ 6 \\ 7 \\ 8 \end{matrix}$$



$$[K] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 \\ 1/2 & -1/2 & 0 & 0 & 0 & 0 & -1/2 & -1/2 \\ -1/2 & 1/2 & 0 & 0 & 0 & 0 & 1/2 & -1/2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 2\sqrt{2} & 0 & 0 & 0 & -2\sqrt{2} \\ 0 & 0 & 0 & 0 & 1/2 & 1/2 & -1/2 & -1/2 \\ 0 & 0 & 0 & 0 & 1/2 & 1/2 & -1/2 & -1/2 \\ -1/2 & 1/2 & 0 & 0 & -1/2 & -1/2 & 1/2+0+1/2 & -1/2+0+1/2 \\ 1/2 & -1/2 & 0 & -2\sqrt{2} & -1/2 & -1/2 & -1/2+1/2 & 1/2+2\sqrt{2}+1/2 \end{bmatrix} \begin{matrix} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \\ 8 \end{matrix}$$

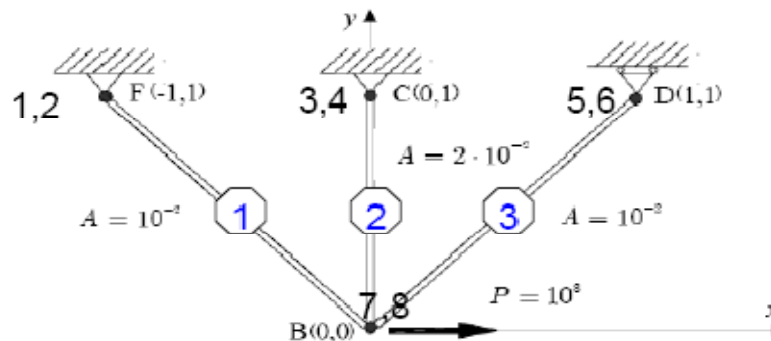
# A truss example: Assembly



$$[K] = \frac{EA}{\sqrt{2}} \begin{bmatrix} 1/2 & -1/2 & 0 & 0 & 0 & 0 & -1/2 & -1/2 \\ -1/2 & 1/2 & 0 & 0 & 0 & 0 & 1/2 & -1/2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 2\sqrt{2} & 0 & 0 & 0 & -2\sqrt{2} \\ 0 & 0 & 0 & 0 & 1/2 & 1/2 & -1/2 & -1/2 \\ 0 & 0 & 0 & 0 & 1/2 & 1/2 & -1/2 & -1/2 \\ -1/2 & 1/2 & 0 & 0 & -1/2 & -1/2 & 1 & 0 \\ 1/2 & -1/2 & 0 & -2\sqrt{2} & -1/2 & -1/2 & 0 & 1+2\sqrt{2} \end{bmatrix} \begin{matrix} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \\ 8 \end{matrix}$$

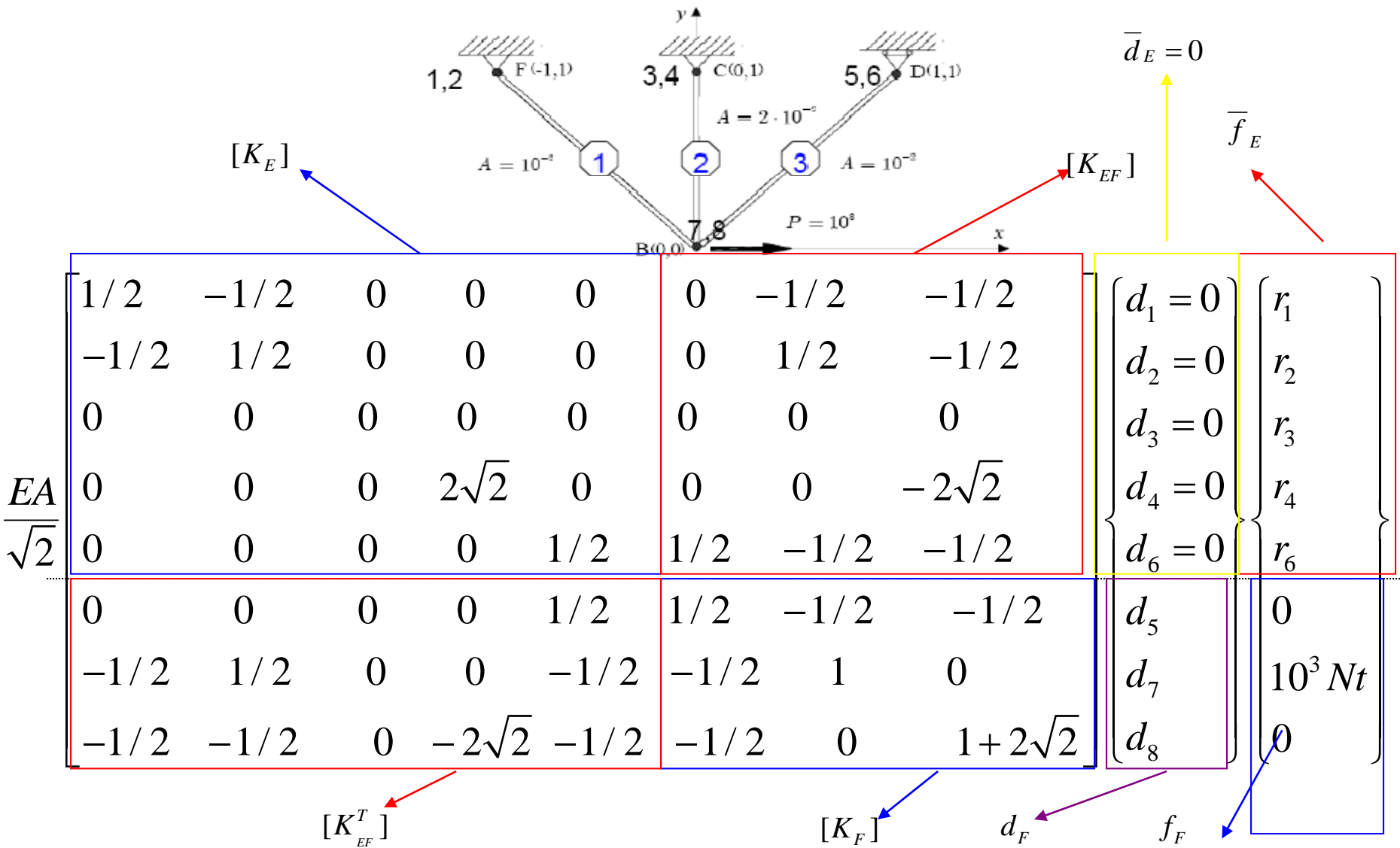


# A truss example: Partition and BCs

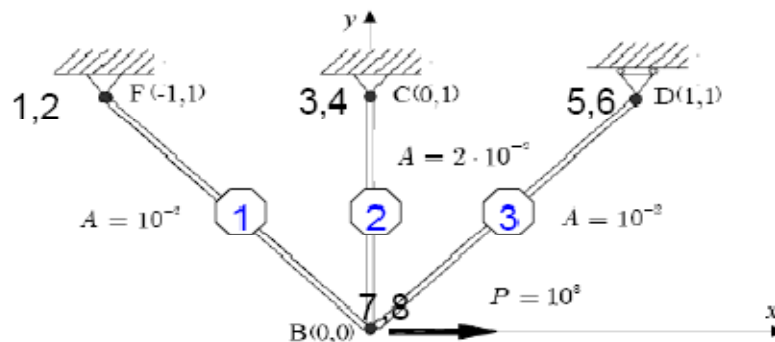


$$EA \begin{bmatrix} 1/2 & -1/2 & 0 & 0 & 0 & 0 & -1/2 & -1/2 \\ -1/2 & 1/2 & 0 & 0 & 0 & 0 & 1/2 & -1/2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 2\sqrt{2} & 0 & 0 & 0 & -2\sqrt{2} \\ \sqrt{2} & 0 & 0 & 0 & 1/2 & 1/2 & -1/2 & -1/2 \\ 0 & 0 & 0 & 0 & 1/2 & 1/2 & -1/2 & -1/2 \\ -1/2 & 1/2 & 0 & 0 & -1/2 & -1/2 & 1 & 0 \\ 1/2 & -1/2 & 0 & -2\sqrt{2} & -1/2 & -1/2 & 0 & 1+2\sqrt{2} \end{bmatrix} \begin{bmatrix} d_1 = 0 \\ d_2 = 0 \\ d_3 = 0 \\ d_4 = 0 \\ d_5 \\ d_6 = 0 \\ d_7 \\ d_8 \end{bmatrix} \begin{bmatrix} r_1 \\ r_2 \\ r_3 \\ r_4 \\ 0 \\ r_6 \\ 10^3 Nt \\ 0 \end{bmatrix}$$

# A truss example: Partition and BCs



# A truss example: Partition and BCs

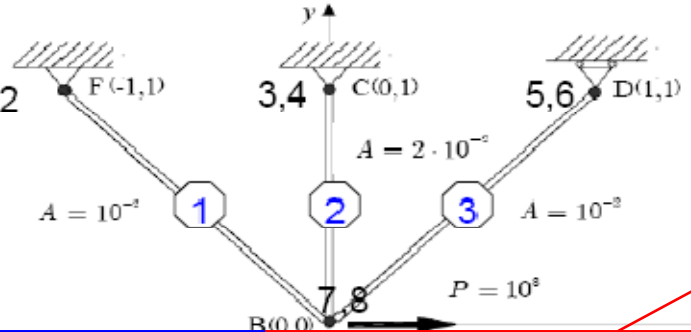


$$\underbrace{\frac{EA}{\sqrt{2}} \begin{bmatrix} 1/2 & -1/2 & -1/2 \\ -1/2 & 1 & 0 \\ -1/2 & 0 & 1+2\sqrt{2} \end{bmatrix}}_{K_F} \underbrace{\begin{Bmatrix} d_5 \\ d_7 \\ d_8 \end{Bmatrix}}_{d_F} = \underbrace{\begin{Bmatrix} 0 \\ 10^3 \text{ Nt} \\ 0 \end{Bmatrix}}_{f_F}$$

$$\begin{Bmatrix} d_5 \\ d_7 \\ d_8 \end{Bmatrix} = \begin{Bmatrix} 0.038284\text{m} \\ 0.033284\text{m} \\ 0.005\text{m} \end{Bmatrix}$$

# A truss example: Reaction calculation

$$\bar{f}_E = K_E \bar{d}_E + K_{EF} d_{F,1,2}$$



$$\bar{d}_E = 0$$

$$\bar{f}_E$$

$$[K_E]$$

$$[K_{EF}]$$

$\frac{EA}{\sqrt{2}}$	1/2	-1/2	0	0	0	0	-1/2	-1/2	$\left. \begin{array}{l} d_1 = 0 \\ d_2 = 0 \\ d_3 = 0 \\ d_4 = 0 \\ d_6 = 0 \end{array} \right\} \begin{array}{l} r_1 \\ r_2 \\ r_3 \\ r_4 \\ r_6 \end{array}$
	-1/2	1/2	0	0	0	0	1/2	-1/2	
	0	0	0	0	0	0	0	0	
	0	0	0	$2\sqrt{2}$	0	0	0	$-2\sqrt{2}$	
	0	0	0	0	1/2	1/2	-1/2	-1/2	
	0	0	0	0	1/2	1/2	-1/2	-1/2	
	-1/2	1/2	0	0	-1/2	-1/2	1	0	$\left. \begin{array}{l} d_5 \\ d_7 \\ d_8 \end{array} \right\} \begin{array}{l} 0 \\ 10^3 Nt \\ 0 \end{array}$
	1/2	-1/2	0	$-2\sqrt{2}$	-1/2	-1/2	0	$1+2\sqrt{2}$	

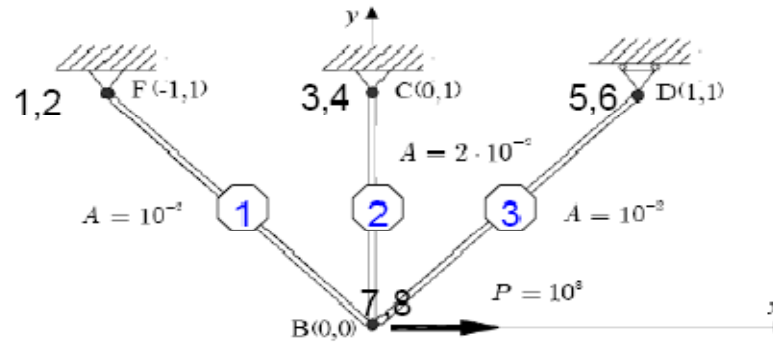
$$[K_{EF}^T]$$

$$[K_F]$$

$$d_F$$

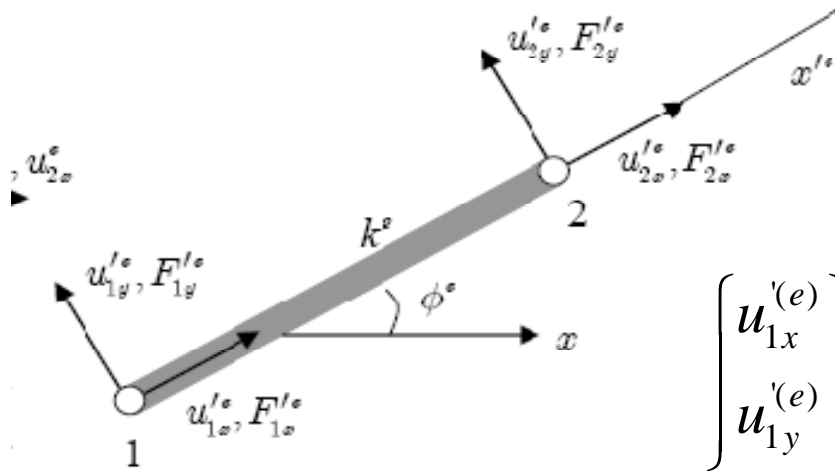
$$f_F$$

# A truss example: Reaction calculation



$$\underbrace{\begin{Bmatrix} r_1 \\ r_2 \\ r_3 \\ r_4 \\ r_6 \end{Bmatrix}}_{f_E} = \frac{EA}{\sqrt{2}} \underbrace{\begin{bmatrix} 0 & -1/2 & 1/2 \\ 0 & 1/2 & -1/2 \\ 0 & 0 & 0 \\ 0 & 0 & -2\sqrt{2} \\ 1/2 & -1/2 & -1/2 \end{bmatrix}}_{K_{EF}} \underbrace{\begin{Bmatrix} d_5 \\ d_7 \\ d_8 \end{Bmatrix}}_{d_F} \Rightarrow \begin{Bmatrix} r_1 \\ r_2 \\ r_3 \\ r_4 \\ r_6 \end{Bmatrix} \equiv \begin{Bmatrix} R_{Fx} \\ R_{Fy} \\ R_{Cx} \\ R_{Cy} \\ R_{Dy} \end{Bmatrix} = \begin{Bmatrix} -1000Nt \\ 1000Nt \\ 0 \\ -1000Nt \\ 0 \end{Bmatrix}$$

# A truss example: Compute the stresses



$$\varepsilon^e = \frac{u_{2x}'^e - u_{1x}'^e}{L^e} \Rightarrow \sigma^e = E^e \frac{u_{2x}'^e - u_{1x}'^e}{L^e}$$

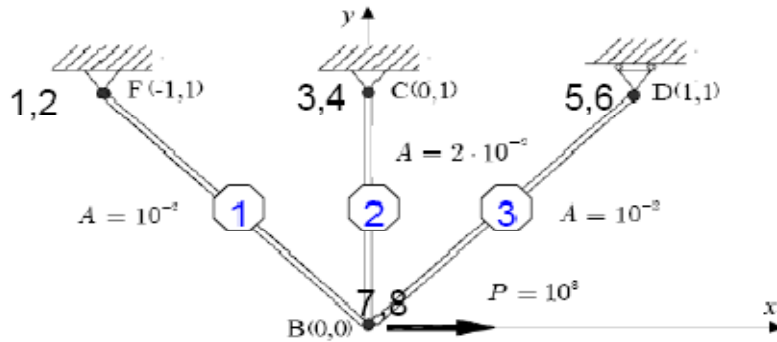
However:

$$\begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix} = \begin{bmatrix} \cos \phi^e & \sin \phi^e & 0 & 0 \\ -\sin \phi^e & \cos \phi^e & 0 & 0 \\ 0 & 0 & \cos \phi^e & \sin \phi^e \\ 0 & 0 & -\sin \phi^e & \cos \phi^e \end{bmatrix} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix}$$

Combining the 2 Eqs gives:

$$\sigma^e = \frac{E^e}{L^e} \begin{bmatrix} -\cos \phi^e & -\sin \phi^e & \cos \phi^e & \sin \phi^e \end{bmatrix} \{d^e\}$$

# A truss example: Compute the stresses



$$\sigma^e = \frac{E^e}{L^e} \begin{bmatrix} -\cos \phi^e & -\sin \phi^e & \cos \phi^e & \sin \phi^e \end{bmatrix} \{d^e\}$$

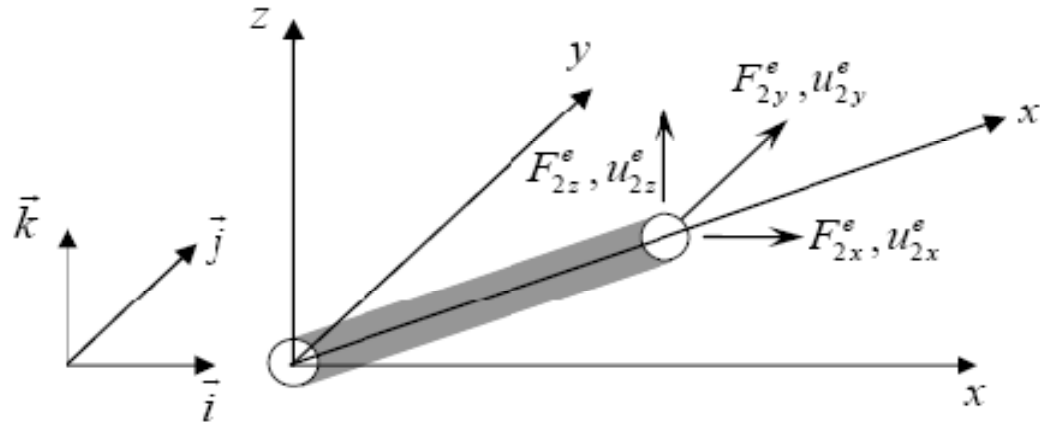
- Applying this to each element, we have:

$$\sigma^{(1)} = \frac{E}{\sqrt{2}} \begin{bmatrix} \frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix} \begin{Bmatrix} 0 \\ 0 \\ 0.038284m \\ 0.005m \end{Bmatrix} = -141.421kPa$$

$$\sigma^{(2)} = E \begin{bmatrix} 0 & -1 & 0 & 1 \end{bmatrix} \begin{Bmatrix} 0 \\ 0 \\ 0.033284m \\ 0.005m \end{Bmatrix} = 50kPa$$

$$\sigma^{(3)} = \frac{E}{\sqrt{2}} \begin{bmatrix} -\frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix} \begin{Bmatrix} 0.038284m \\ 0 \\ 0.033284m \\ 0.005m \end{Bmatrix} = 0kPa$$

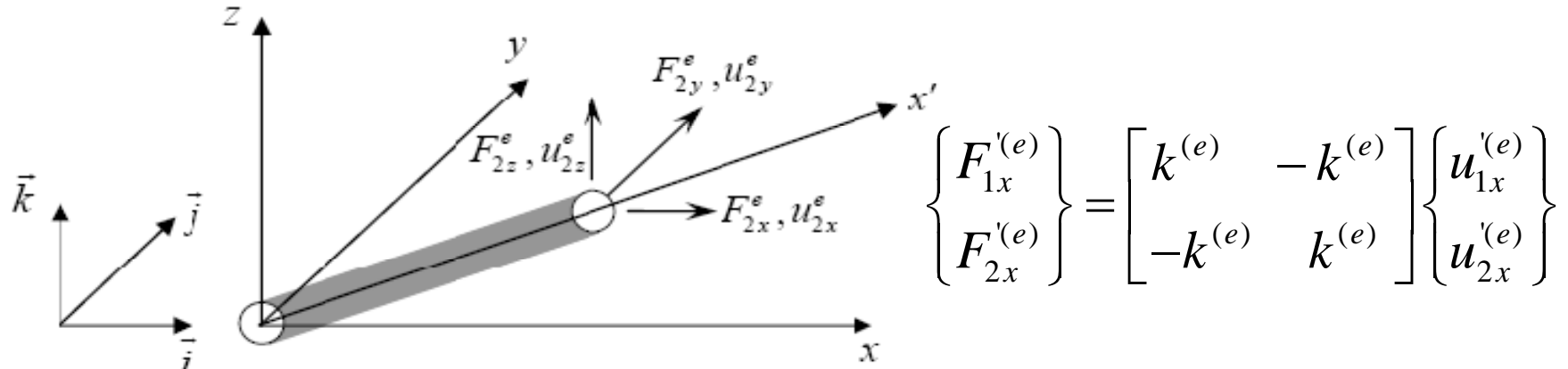
# Three-dimensional (space) truss structures



- The local stiffness eqs are exactly as before, i.e. 
$$\begin{Bmatrix} F_{1x}^{(e)} \\ F_{2x}^{(e)} \end{Bmatrix} = \begin{bmatrix} k^{(e)} & -k^{(e)} \\ -k^{(e)} & k^{(e)} \end{bmatrix} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{2x}^{(e)} \end{Bmatrix}, \quad k^{(e)} = A^e E^e / L^e$$
- However, we now have to include displacements and forces in the  $x$ ,  $y$  and  $z$  axes.



# Three-dimensional (space) truss structures



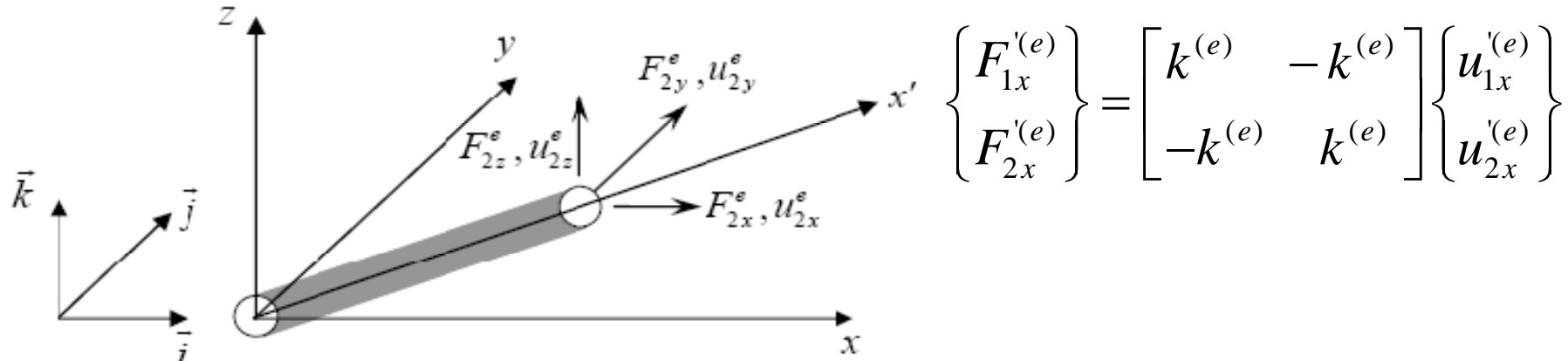
- A **unit vector** along the direction  $x'$  of a 3D truss element has the components (**direction cosines of the axes between  $x'$  and  $x, y, z$ , respectively**):

$$l_s^e = \frac{x_2^e - x_1^e}{L^e}, \quad m_s^e = \frac{y_2^e - y_1^e}{L^e}, \quad n_s^e = \frac{z_2^e - z_1^e}{L^e}$$

$$L^e = \sqrt{(x_2^e - x_1^e)^2 + (y_2^e - y_1^e)^2 + (z_2^e - z_1^e)^2}$$

where  $(x_1^e, y_1^e, z_1^e)$  and  $(x_2^e, y_2^e, z_2^e)$  are the nodal coordinates in the  $(x, y, z)$  system.

# Three-dimensional (space) truss structures



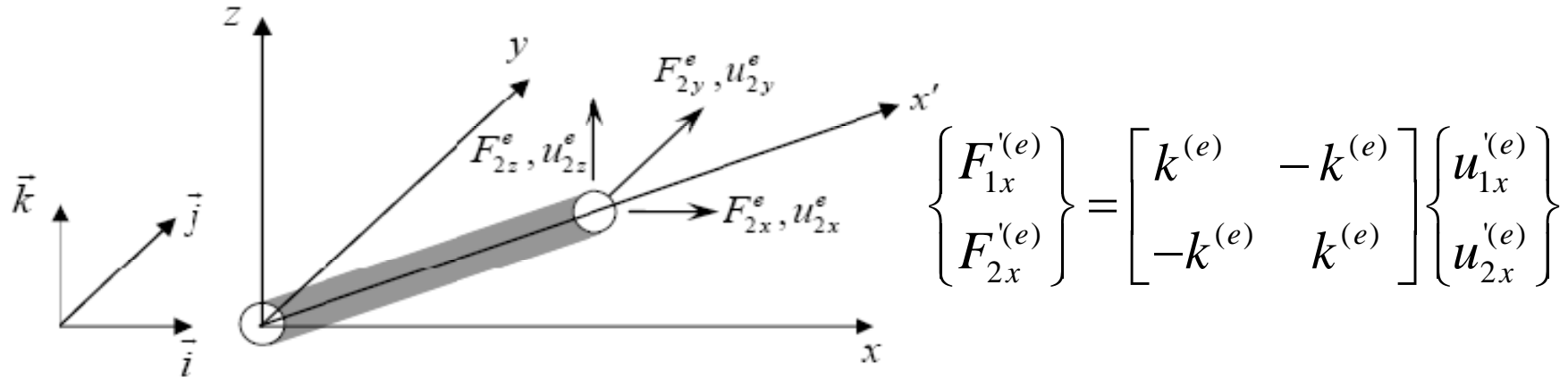
- The displacement transformation then takes the form:

$$\begin{Bmatrix} u_{1x}'^{(e)} \\ u_{2x}'^{(e)} \end{Bmatrix} = \underbrace{\begin{bmatrix} l_s^e & m_s^e & n_s^e & 0 & 0 & 0 \\ 0 & 0 & 0 & l_s^e & m_s^e & n_s^e \end{bmatrix}}_{[T^e]} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{1z}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \\ u_{2z}^{(e)} \end{Bmatrix} \equiv [T^e] \{d^e\}$$

$$\begin{Bmatrix} F_{1x}'^{(e)} \\ F_{2x}'^{(e)} \end{Bmatrix} = [T^e] \{F^e\}$$

- Similar transformation is applied for the forces:

# Three-dimensional (space) truss structures



- Similarly to the derivation for 2D trusses, **the stiffness matrix in global coordinates** is then:

$$\underbrace{[K]}_{6 \times 6} \equiv \underbrace{[T^e]^T}_{6 \times 2} \underbrace{\{K^e\}}_{2 \times 2} \underbrace{[T^e]}_{2 \times 6}$$

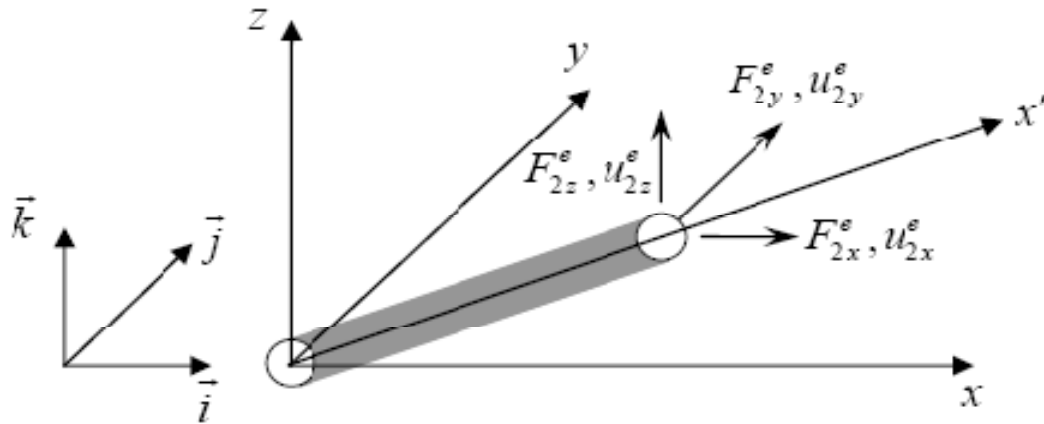
# Stiffness of a space element

$$[K^{(e)}] = \frac{E^e A^e}{L^e} \begin{bmatrix} l_s^{e2} & m_s^e l_s^e & n_s^e l_s^e & -l_s^{e2} & -m_s^e l_s^e & -n_s^e l_s^e \\ m_s^e l_s^e & m_s^{e2} & m_s^e n_s^e & -m_s^e l_s^e & -m_s^{e2} & -m_s^e n_s^e \\ n_s^e l_s^e & m_s^e n_s^e & n_s^{e2} & -n_s^e l_s^e & -m_s^e n_s^e & -n_s^{e2} \\ -l_s^{e2} & -m_s^e l_s^e & -n_s^e l_s^e & l_s^{e2} & m_s^e l_s^e & n_s^e l_s^e \\ -m_s^e l_s^e & -m_s^{e2} & -m_s^e n_s^{e2} & m_s^e l_s^e & m_s^{e2} & m_s^e n_s^e \\ -n_s^e l_s^e & -m_s^e n_s^e & -n_s^{e2} & n_s^e l_s^e & m_s^e n_s^e & n_s^{e2} \end{bmatrix}$$

where:  $l_s^e = \frac{x_2^e - x_1^e}{L^e}, \quad m_s^e = \frac{y_2^e - y_1^e}{L^e}, \quad n_s^e = \frac{z_2^e - z_1^e}{L^e}$

$$L^e = \sqrt{(x_2^e - x_1^e)^2 + (y_2^e - y_1^e)^2 + (z_2^e - z_1^e)^2}$$

# Computing the stresses in a space truss element



$$\epsilon^e = \frac{u_{2x}^e - u_{1x}^e}{L^e} \Rightarrow$$

$$\sigma^e = E^e \frac{u_{2x}^e - u_{1x}^e}{L^e} = \frac{E^e}{L^e} (u_{2x}^e - u_{1x}^e)$$

➤ However:

$$\begin{Bmatrix} u_{1x}^{(e)} \\ u_{2x}^{(e)} \end{Bmatrix} = \begin{bmatrix} l_s^e & m_s^e & n_s^e & 0 & 0 & 0 \\ 0 & 0 & 0 & l_s^e & m_s^e & n_s^e \end{bmatrix} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{1z}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \\ u_{2z}^{(e)} \end{Bmatrix} \equiv [T^e] \{d^e\}$$

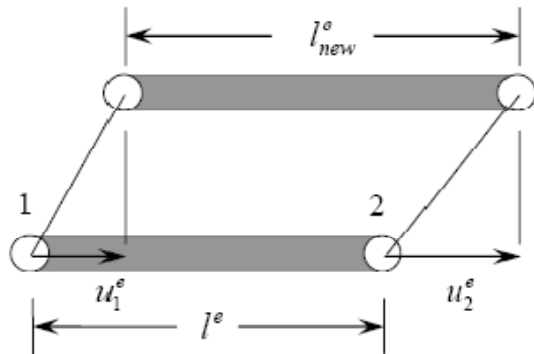
➤ Combining the 2 Eqs gives:

$$\sigma^e = \frac{E^e}{L^e} \begin{bmatrix} -l_s^e & -m_s^e & -n_s^e & l_s^e & m_s^e & n_s^e \end{bmatrix} \{d^e\}$$

# Accounting for thermal effects in truss analysis

- Consider a truss structure that is heated. We need to account for thermal expansion effects. Note:  $\underbrace{\varepsilon^e}_{\text{total strain}} = \varepsilon_{\text{elastic}}^e + \varepsilon_{\text{thermal}}^e$
- Hooke's law is now modified as follows (using the x' coordinate system):

$$\sigma^e = E^e \varepsilon_{\text{elastic}}^e = E^e \left( \underbrace{\varepsilon^e}_{\text{total strain}} - \varepsilon_{\text{thermal}}^e \right) = E^e \left( \frac{u_2^e - u_1^e}{L^e} - \alpha^e \Delta T^e \right)$$



$$F_2^{(e)} = -F_1^{(e)} = p^e = A^e \sigma^e = A^e E^e \left( \frac{u_2^e - u_1^e}{L^e} - \underbrace{\alpha^e \Delta T^e}_{\varepsilon_0^e} \right) =$$

$$= k^e (u_2^e - u_1^e) - A^e E^e \varepsilon_0^e, \quad k^e = \frac{A^e E^e}{L^e}$$

$$\begin{Bmatrix} F_1^{(e)} \\ F_2^{(e)} \end{Bmatrix} + \underbrace{A^e E^e \varepsilon_0^e}_{\text{Thermal nodal vector}} \begin{Bmatrix} -1 \\ 1 \end{Bmatrix} = \begin{bmatrix} k^{(e)} & -k^{(e)} \\ -k^{(e)} & k^{(e)} \end{bmatrix} \begin{Bmatrix} u_1^{(e)} \\ u_2^{(e)} \end{Bmatrix}$$

# Element equations with thermal effects

- Expanding these equations to include nodal displacements in the  $y'$  axis gives:

$$\underbrace{\begin{Bmatrix} F_{1x}^{(e)} \\ F_{1y}^{(e)} \\ F_{2x}^{(e)} \\ F_{2y}^{(e)} \end{Bmatrix}}_{\{F'^{(e)}\}} + \underbrace{A^e E^e \varepsilon_0}_{\{F_{thermal}^e\}} \underbrace{\begin{Bmatrix} -1 \\ 0 \\ 1 \\ 0 \end{Bmatrix}}_{\{F_{thermal}^e\}} = \underbrace{k^{(e)}}_{[K^{(e)}]} \underbrace{\begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix}}_{\{d'^{(e)}\}}$$

- We need to transform this to  $x$  and  $y$  displacements (our degrees of freedom for this element)

$$\{d^e\} = [T^e] \{d'^e\}$$

$$\{F'^e\} = [T^e] \{F^e\}$$

# Element equations with thermal effects

$$\underbrace{\begin{Bmatrix} F_{1x}^{(e)} \\ F_{1y}^{(e)} \\ F_{2x}^{(e)} \\ F_{2y}^{(e)} \end{Bmatrix}}_{\{F^{(e)}\}} + \underbrace{A^e E^e \varepsilon_0^e \begin{Bmatrix} -1 \\ 0 \\ 1 \\ 0 \end{Bmatrix}}_{\{F_{thermal}^{(e)}\}} = \underbrace{k^{(e)} \begin{bmatrix} 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}}_{[K^{(e)}]} \underbrace{\begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix}}_{\{d^{(e)}\}}$$

$$\{d^{(e)}\} = [T^e] \{d^e\}$$

$$\{F^{(e)}\} = [T^e] \{F^e\}$$

- We can transform these element equations as follows:

$$[T^e] \{F^e\} + \{F_{thermal}^{(e)}\} = [K^{(e)}] [T^e] \{d^e\}$$

- From these equations, we conclude that:

$$\underbrace{\{F^e\} + [T^e]^T \{F_{thermal}^{(e)}\}}_{\{F_{thermal}^e\}} = \underbrace{[T^e]^T [K^{(e)}] [T^e]}_{[K^{(e)}]} \{d^e\}$$



# Element equations with thermal effects

$$\underbrace{\{F^e\} + [T^e]^T \{F_{thermal}^e\}}_{\{F_{thermal}^e\}} = \underbrace{[T^e]^T [K^{(e)}] [T^e]}_{[K^{(e)}]} \{d^e\}$$

Use:  $[T^{(e)}] = \begin{bmatrix} \cos \phi^e & \sin \phi^e & 0 & 0 \\ -\sin \phi^e & \cos \phi^e & 0 & 0 \\ 0 & 0 & \cos \phi^e & \sin \phi^e \\ 0 & 0 & -\sin \phi^e & \cos \phi^e \end{bmatrix}$  (for 2D trusses)

Finally we obtain:

$$\underbrace{\begin{Bmatrix} F_{1x}^{(e)} \\ F_{1y}^{(e)} \\ F_{2x}^{(e)} \\ F_{2y}^{(e)} \end{Bmatrix}}_{\{F^{(e)}\}} + A^e E^e \epsilon_0^e \underbrace{\begin{Bmatrix} -\cos \phi^e \\ -\sin \phi^e \\ \cos \phi^e \\ \sin \phi^e \end{Bmatrix}}_{\{F_{thermal}^e\}} = k^{(e)} \underbrace{\begin{bmatrix} \cos^2 \phi^e & \sin \phi^e \cos \phi^e & -\cos^2 \phi^e & -\sin \phi^e \cos \phi^e \\ \sin \phi^e \cos \phi^e & \sin^2 \phi^e & -\sin \phi^e \cos \phi^e & -\sin^2 \phi^e \\ -\cos^2 \phi^e & -\sin \phi^e \cos \phi^e & \cos^2 \phi^e & \sin \phi^e \cos \phi^e \\ -\sin \phi^e \cos \phi^e & -\sin^2 \phi^e & \sin \phi^e \cos \phi^e & \sin^2 \phi^e \end{bmatrix}}_{[K^{(e)}]} \underbrace{\begin{Bmatrix} u_{1x}^{(e)} \\ u_{1y}^{(e)} \\ u_{2x}^{(e)} \\ u_{2y}^{(e)} \end{Bmatrix}}_{d^{(e)}}$$

# Element equations with thermal effects

- What do you need to do to account for thermal effects in truss analysis?
- For each truss element that is heated, simply **add to the element force, the following extra term**

$$\underbrace{A^e E^e \varepsilon_0^e \begin{Bmatrix} -\cos \phi^e \\ -\sin \phi^e \\ \cos \phi^e \\ \sin \phi^e \end{Bmatrix}}_{\{F_{thermal}^e\}}, \text{ where } \varepsilon_0^e = \alpha^e \Delta T^e$$

- You will need to define at which truss elements thermal effects take place and for each of them read the values  $\alpha^e$  and  $\Delta T^e$ .

# Principle of minimum potential energy

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- An alternative **equivalent approach** to solving many structural problems is the **principle of minimum potential energy**.

*From all the possible compatible displacements of a structure, the one that minimizes the total potential energy is the exact solution.*

Potential energy  
for given  
displacements

=

Strain energy  
for these given  
displacements

-

Work done by external  
loads on these given  
displacements

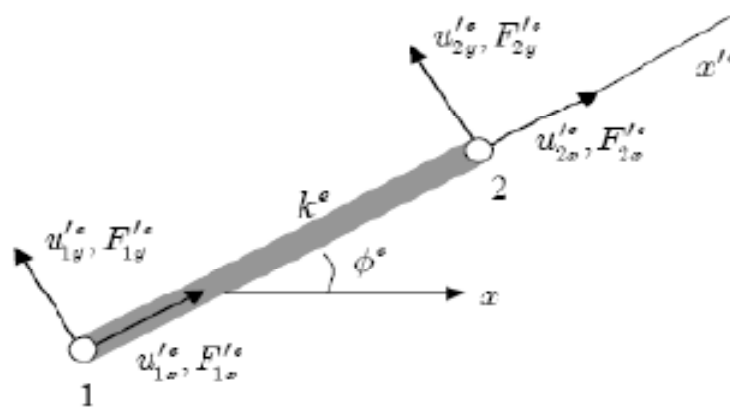
# Principle of minimum potential energy

- Let us see this principle applied to the truss problems discussed earlier.

Assembly process

$$\min_{\{d\}} \sum_e PE^e, PE^e = U^e - W^e$$

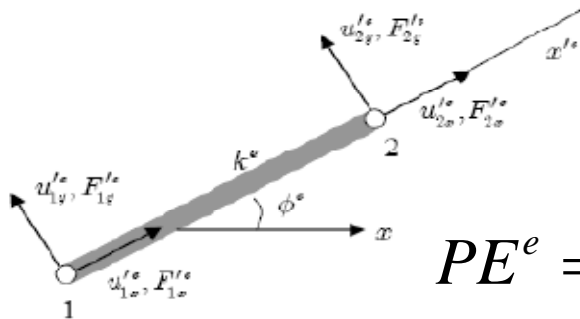
$$PE^e = \int_{\Omega^e} \underbrace{\frac{1}{2} \sigma^e \varepsilon^e}_{\text{Elastic strain energy density (work/volume)}} dV^e - \underbrace{(F_{1x}^{(e)} u_{1x}^{(e)} + F_{2x}^{(e)} u_{2x}^{(e)})}_{\text{External Work}}$$



# Principle of minimum potential energy

- Lets apply this principle to one truss element. We need to minimize with respect to the nodal displacements (local coordinates)  $u_{1x}^{(e)}, u_{2x}^{(e)}$ . Recall that:

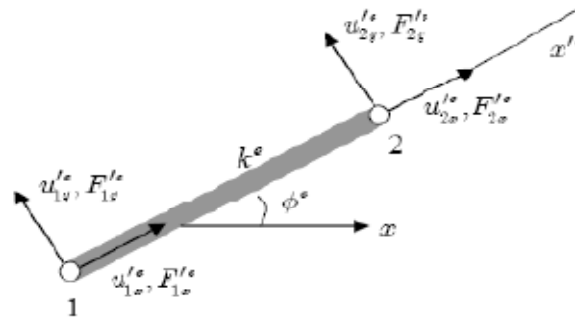
$$\varepsilon^e = \frac{u_{2x}^{(e)} - u_{1x}^{(e)}}{L^e}, \sigma^e = E^e \varepsilon^e$$



$$PE^e = \int_{\Omega^e} \frac{1}{2} E^e \varepsilon^{e2} dV^e - F_{1x}^{(e)} u_{1x}^{(e)} - F_{2x}^{(e)} u_{2x}^{(e)} =$$

$$\int_{\Omega^e} \frac{1}{2} E^e \left( \frac{u_{2x}^{(e)} - u_{1x}^{(e)}}{L^e} \right)^2 \underbrace{dV^e}_{A^e dx'} - F_{1x}^{(e)} u_{1x}^{(e)} - F_{2x}^{(e)} u_{2x}^{(e)}$$

# Principle of minimum potential energy



$$\min_{u_{1x}^{(e)}, u_{2x}^{(e)}} PE^e = \min_{u_{1x}^{(e)}, u_{2x}^{(e)}} E^e A^e L^e \frac{1}{2} \left( \frac{u_{2x}^{(e)} - u_{1x}^{(e)}}{L^e} \right)^2 - F_{1x}^{(e)} u_{1x}^{(e)} - F_{2x}^{(e)} u_{2x}^{(e)}$$

- Take partial derivatives of  $PE^e$  wrt  $u_{1x}^{(e)}, u_{2x}^{(e)}$  :

$$\begin{aligned} \frac{\partial PE^e}{\partial u_{1x}^{(e)}} = 0 &\Rightarrow \frac{E^e A^e}{L^e} (u_{1x}^{(e)} - u_{2x}^{(e)}) - F_{1x}^{(e)} = 0 \\ \frac{\partial PE^e}{\partial u_{2x}^{(e)}} = 0 &\Rightarrow \frac{E^e A^e}{L^e} (u_{2x}^{(e)} - u_{1x}^{(e)}) - F_{2x}^{(e)} = 0 \end{aligned} \Rightarrow \begin{Bmatrix} F_{1x}^{(e)} \\ F_{2x}^{(e)} \end{Bmatrix} = \begin{bmatrix} k^{(e)} & -k^{(e)} \\ -k^{(e)} & k^{(e)} \end{bmatrix} \begin{Bmatrix} u_{1x}^{(e)} \\ u_{2x}^{(e)} \end{Bmatrix}$$

These are the same Eqs as those obtained with the direct method!

# Principle of minimum potential energy

- In general (not just for mechanics problems!), the principle of minimum potential energy takes the following form:

$$\min_{\{d\}} \sum_e PE^e = \min_{\{d\}} \left( \sum_e \left( \frac{1}{2} \{d^{(e)}\}^T [K^{(e)}] \{d^{(e)}\} \right) - \{d\}^T \{F\} \right)$$

or after assembly:

$$\min_{\{d\}} PE = \min_{\{d\}} \left( \frac{1}{2} \{d\}^T [K] \{d\} - \{d\}^T \{F\} \right)$$

- Note that the minimization gives us the familiar solution:  $[K]\{d\} = \{F\}$

# Principle of minimum potential energy

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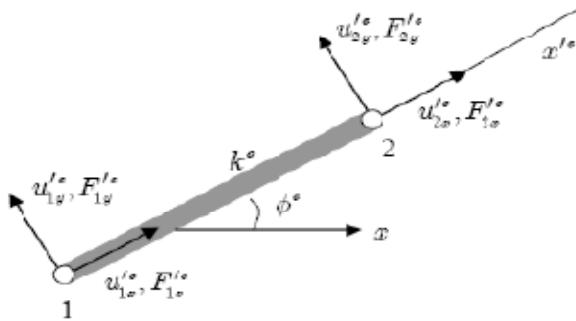
- We will not use this method to repeat the truss calculations.
- However, it will be our starting point for computing **the stiffness of beam elements (lecture 4)**.
- The method of minimum potential energy can be applied to many problems not related to mechanics – however there are many problems where this technique is not applicable.
- After discussing beam bending problems, we will need to look for more powerful (‘unfortunately’ also more mathematical) methods (weak (Galerkin) formulations).



# Revisiting the 2-node truss element

- Up to now we used the direct method to express the nodal loads vs. nodal displacements for the 2-node truss element.
- Let us **linearly interpolate** the displacement  $u_x^{(e)}$  of any point  $x'$  in the element in terms of the nodal displacements:

$$u_x^{(e)} = \left(1 - \frac{x'}{L^e}\right)u_{1x}^{(e)} + \frac{x'}{L^e}u_{2x}^{(e)} = \underbrace{\left[1 - \frac{x'}{L^e}, \frac{x'}{L^e}\right]}_{\substack{N_1^{(e)} \quad N_2^{(e)} \\ \text{element basis} \\ \text{functions}}} \underbrace{\begin{Bmatrix} u_{1x}^{(e)} \\ u_{2x}^{(e)} \end{Bmatrix}}_{\substack{\text{Nodal} \\ \text{displacements}}} = \underbrace{[N^{(e)}]}_{\substack{\text{basis functions} \\ \text{matrix}}} \{d^{(e)}\}$$



- The strain in this 2-node element  $\varepsilon^e = \frac{du_x^{(e)}}{dx'}$  can now be computed as follows:

$$u_x^{(e)} = [N^{(e)}]\{d^{(e)}\} \Rightarrow \varepsilon^e = \frac{du_x^{(e)}}{dx'} = \frac{d[N^{(e)}]\{d^{(e)}\}}{dx'} = \left[\frac{dN^{(e)}}{dx'}\right]\{d^{(e)}\} \equiv [B^{(e)}]\{d^{(e)}\}$$

$$\underbrace{[B^{(e)}]}_{\substack{\text{Gradient of} \\ \text{basis functions} \\ \text{matrix}}} = \left[\frac{dN_1^{(e)}}{dx'}, \frac{dN_2^{(e)}}{dx'}\right] = \left[-\frac{1}{L^e}, \frac{1}{L^e}\right] \Rightarrow \varepsilon^e = \left[-\frac{1}{L^e}, \frac{1}{L^e}\right] \begin{Bmatrix} u_{1x}^{(e)} \\ u_{2x}^{(e)} \end{Bmatrix} = \frac{u_{2x}^{(e)} - u_{1x}^{(e)}}{L^e}$$

- This is exactly the same approximation we used before (**constant strain element**)

# Revisiting the 2-node truss element

- Let us use these interpolation formulas  $u^{(e)} = [N^{(e)}]\{d^{(e)}\}$ ,  $\varepsilon^e = [B^{(e)}]\{d^{(e)}\}$  to express the potential energy in a format that will become very familiar as we proceed forward in this course.

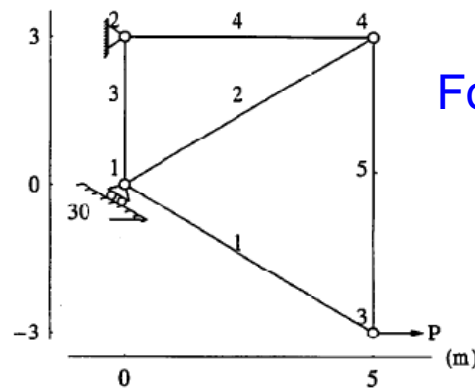
$$\begin{aligned}
 PE^e &= \int_{\Omega^e} \frac{1}{2} E^e \varepsilon^{e2} dV^e - F_{1x}^{(e)} u_{1x}^{(e)} - F_{2x}^{(e)} u_{2x}^{(e)} = \\
 &= \int_{\Omega^e} \frac{1}{2} \underbrace{\{d^{(e)}\}^T}_{\varepsilon^{eT}} [B^e]^T E^e [B^e] \underbrace{\{d^{(e)}\}}_{\varepsilon^e} \underbrace{dV^e}_{A^e dx'} - F_{1x}^{(e)} u_{1x}^{(e)} - F_{2x}^{(e)} u_{2x}^{(e)} = \\
 &= \frac{1}{2} \{d^{(e)}\}^T \underbrace{\left[ \int_{\Omega^e} [B^e]^T E^e [B^e] dV^e \right]}_{\substack{[K^e] \\ \text{Element stiffness} \\ \text{matrix}}} \{d^{(e)}\} - F_{1x}^{(e)} u_{1x}^{(e)} - F_{2x}^{(e)} u_{2x}^{(e)}
 \end{aligned}$$

- For now **these calculations are identical to our earlier approach!** Indeed:

$$[K^e] = \int_{\Omega^e} [B^e]^T E^e [B^e] dV^e = \int_0^{L^e} [B^e]^T E^e [B^e] A^e dx' = \int_0^{L^e} \begin{bmatrix} -\frac{1}{L^e} \\ \frac{1}{L^e} \end{bmatrix} E^e \begin{bmatrix} -\frac{1}{L^e} & \frac{1}{L^e} \end{bmatrix} A^e dx' = \underbrace{\frac{A^e E^e}{L^e}}_{k^e} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$$

# Truss analysis with displacement constraints

- Up to this point, we imposed essential boundary conditions in terms of prescribed  $x$ - or  $y$ - nodal displacements. How about if the support is inclined as in the figure below:



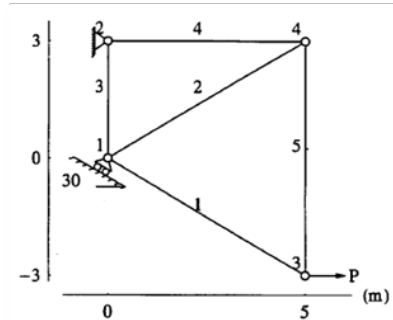
For this problem, the constraint is that there is no normal displacement at the support 1

- Here, we don't know the displacements at node 1 but we know the relation between  $u_{1x}$  and  $u_{1y}$ . In general we write these constraints on our nodal degrees of freedom as:

$$Cd=q.$$

# Truss analysis with displacement constraints

- Note that at node 1 we don't have essential boundary conditions – we have a displacement constraint.



- To solve this problem we use the principle of minimum potential energy with the constraint  $Cd=q$ :

Find  $\mathbf{d}$  and  $\lambda$  such that

$$\min_{\{d\}} L = \underbrace{\frac{1}{2} \{d\}^T [K] \{d\} - \{d\}^T \{F\}}_{\text{Potential energy of unconstrained system}} + \underbrace{\{\lambda\}^T}_{\text{Lagrange multiplier enforcing the constraint}} \underbrace{([C] \{d\} - \{q\})}_{\text{Constraint}}$$

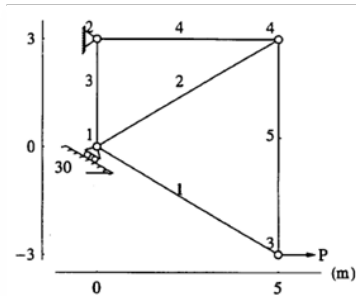
- Here, we enforce the constraint using Lagrange multipliers.

# Truss analysis with displacement constraints

Find  $\mathbf{d}$  and  $\lambda$  such that

$$\min_{\{d\}} L = \underbrace{\frac{1}{2} \{d\}^T [K] \{d\} - \{d\}^T \{F\}}_{\text{Potential energy of unconstrained system}} + \underbrace{\{ \lambda \}^T}_{\text{Lagrange multiplier enforcing the constraint}} \underbrace{([C] \{d\} - \{q\})}_{\text{Constraint}}$$

- $\lambda$  is the Lagrange multiplier that enforces the constraint – it is nothing else but the reaction force at node 1 (normal to the support!)



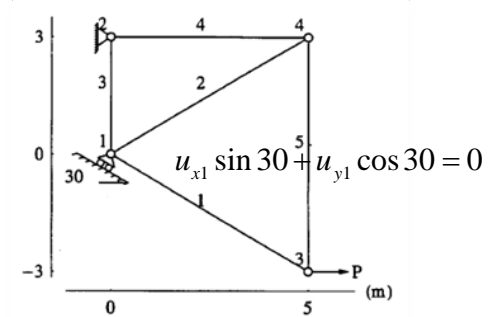
- Minimization is now performed with respect to **both  $\mathbf{d}$  and  $\lambda$** .

$$\begin{pmatrix} K & C^T \\ C & 0 \end{pmatrix} \begin{pmatrix} d \\ \lambda \end{pmatrix} = \begin{pmatrix} F \\ q \end{pmatrix} \Rightarrow$$

Apply essential boundary conditions and then solve for  $\{d_F\}$  and  $\lambda$  (here, reaction force at node 1)

# Displacement constraints: Implementation

- How do you implement this in the MatLab libraries of HW1?



- Introduce the constraints in the InputData.m and then modify the stiffness and load vectors in the NodalSoln.m.
- Apply the essential boundary conditions first before you augment the reduced global equations ( $K_f$ ) with the Lagrange multiplier.

InputData.m

```
C = zeros(1,neq-length(debc));  
  
C(1) = sin(pi/6); C(2) = cos(pi/6);  
q = 0;  
  
% Read information for constraints  
%The dimension of C is neq minus the  
% prescribed DOF via essential BCs  
% Here there is only one constraint
```

# Displacement constraints: NodalSoln.m

```
function [d, rf, lambda] = NodalSoln(K, R, debc, ebcVals, C, q)
```

% K=global stiffness, R=global force, debc=degrees of freedom with specified values, ebcVals=specified displacements

```
dof = length(R); % Extract the total degrees-of-freedom
```

```
df = setdiff(1:dof, debc); % Sets the difference between two sets of indices, i.e. the global degrees of freedom minus the  
% degrees of freedom with prescribed essential boundary conditions
```

```
Kf = K(df, df); % Remove eqs. corresponding to prescribed displacements
```

```
Rf = R(df) - K(df, debc)*ebcVals; % Modify the remaining load vector to account for the essential boundary conditions
```

```
[m n] = size(C); % Extract number of constraints
```

```
Kf = [Kf C; % Augment global equations with the Lagrange multipliers  
C zeros(m)];
```

$$\begin{pmatrix} K & C^T \\ C & 0 \end{pmatrix} \begin{pmatrix} d \\ \lambda \end{pmatrix} = \begin{pmatrix} F \\ q \end{pmatrix}$$

```
Rf = [Rf;q]; % Augment load vector
```

```
dfVals = Kf\Rf; % Solve the linear system of equations. Here for simplicity, we use Gauss elimination.
```

```
d = zeros(dof,1); % Restore the solution vector (i.e. include back the nodes with prescribed displacements).
```

```
d(debc) = ebcVals; % Use the originally established ordering of the nodes.
```

```
d(df) = dfVals(1:(length(dfVals)-m));
```

```
rf = K(debc,:)*d - R(debc); % Calculate the reaction vector at nodes with prescribed displacements
```

```
lambda = dfVals((length(dfVals)-m+1):length(dfVals)); % Calculate Lagrange multipliers (reactions at nodes with constraints)
```