
MAE4700/5700
**Finite Element Analysis for
Mechanical and Aerospace Design**

Cornell University, Fall 2009

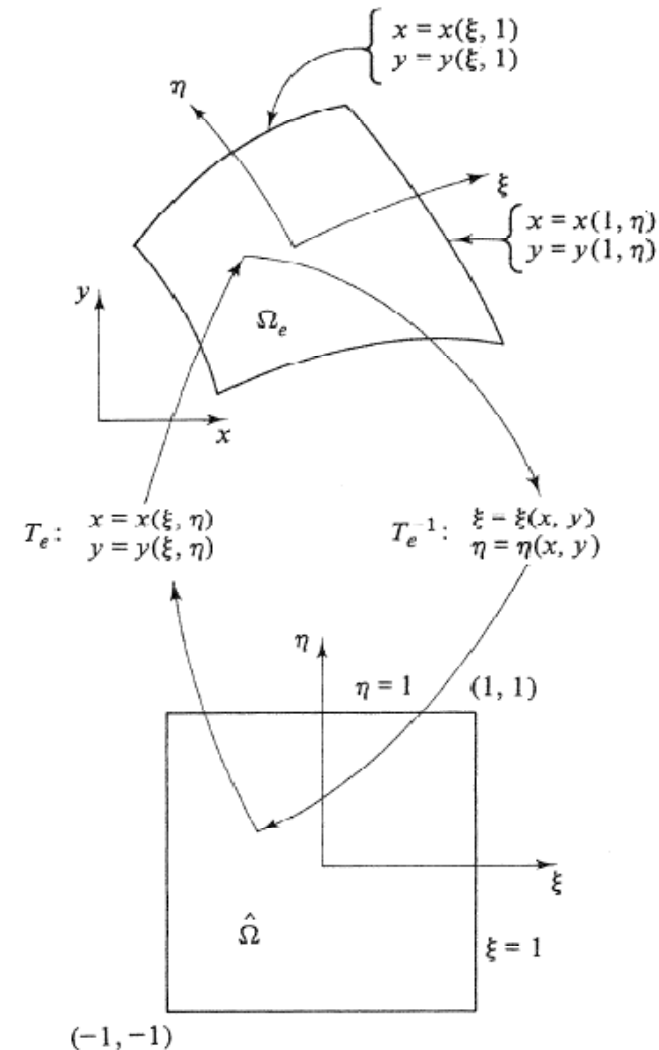
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Two-dimensional element calculations

- Similarly to 1D, we here also define the concept of a **master element in natural coordinates** (ξ, η)
- Performing calculations on elements with curved sides as shown is difficult in the (x, y) coordinates.
 - For example, the limits of integration will vary from element to element

Can we find a master element $\hat{\Omega}$ that is mapped appropriately to each element Ω_e ? Then with transformation $\hat{\Omega} \rightarrow \Omega_e$ we can perform all integrations for all elements at the master element.



Two-dimensional element calculations

- We consider the master element to be a square with: $-1 \leq \xi, \eta \leq 1$
- We define the map $\hat{\Omega} \rightarrow \Omega_e$ as follows

$$x = x(\xi, \eta)$$

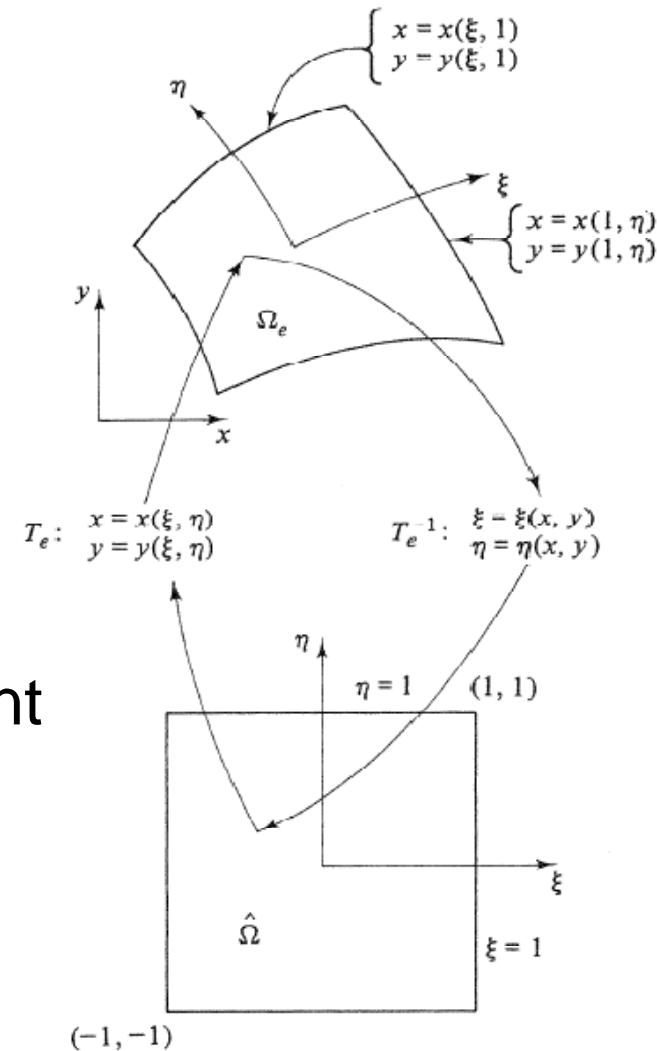
$$T_e :$$

$$y = y(\xi, \eta)$$

- Note that the boundary segment $\xi = 1$ of $\hat{\Omega}$ is mapped on the curved segment $x = x(1, \eta)$

$y = y(1, \eta)$ with η as a parameter.

Similarly for the other 3 boundary segments.

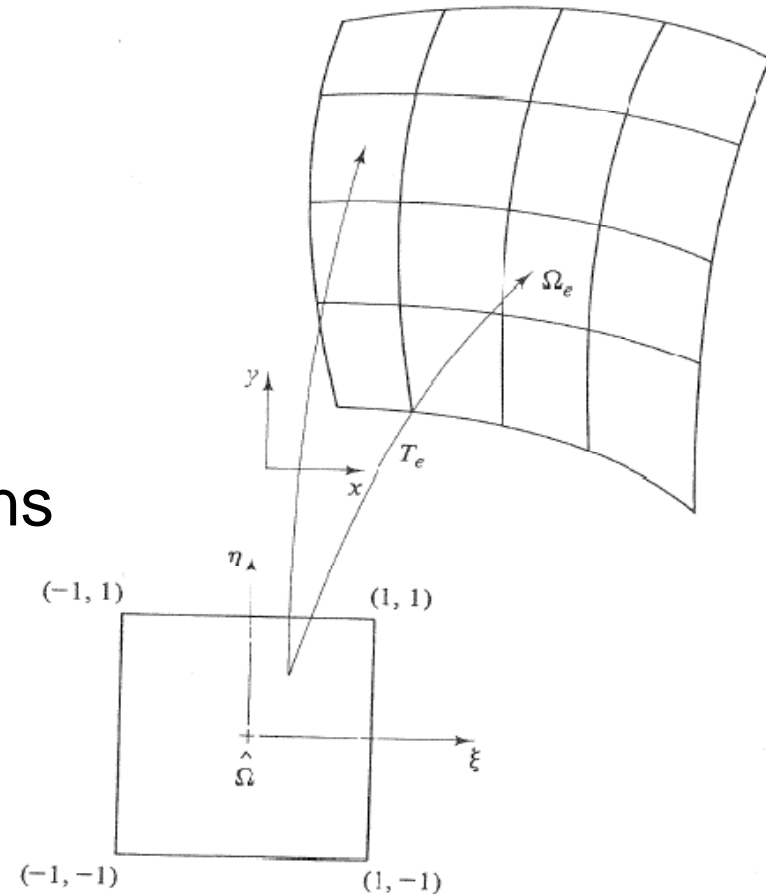


Two-dimensional element calculations

- The complete finite element mesh is seen as a sequence of transformations

$\{T_1, T_2, \dots, T_E\}$ with each of them mapping $T_e: \hat{\Omega} \rightarrow \Omega_e, e=1, 2, \dots, E$

- We need to figure out how we can transfer all needed calculations in each element Ω_e to the master element $\hat{\Omega}$.
- The plan is to make all these calculations defined in terms of the mappings $\{T_1, T_2, \dots, T_E\}$.



Transformation equations

- We assume that the mappings
 $x = x(\xi, \eta)$, $y = y(\xi, \eta)$
are differentiable. We can then write:

$$dx = \frac{\partial x}{\partial \xi} d\xi + \frac{\partial x}{\partial \eta} d\eta,$$

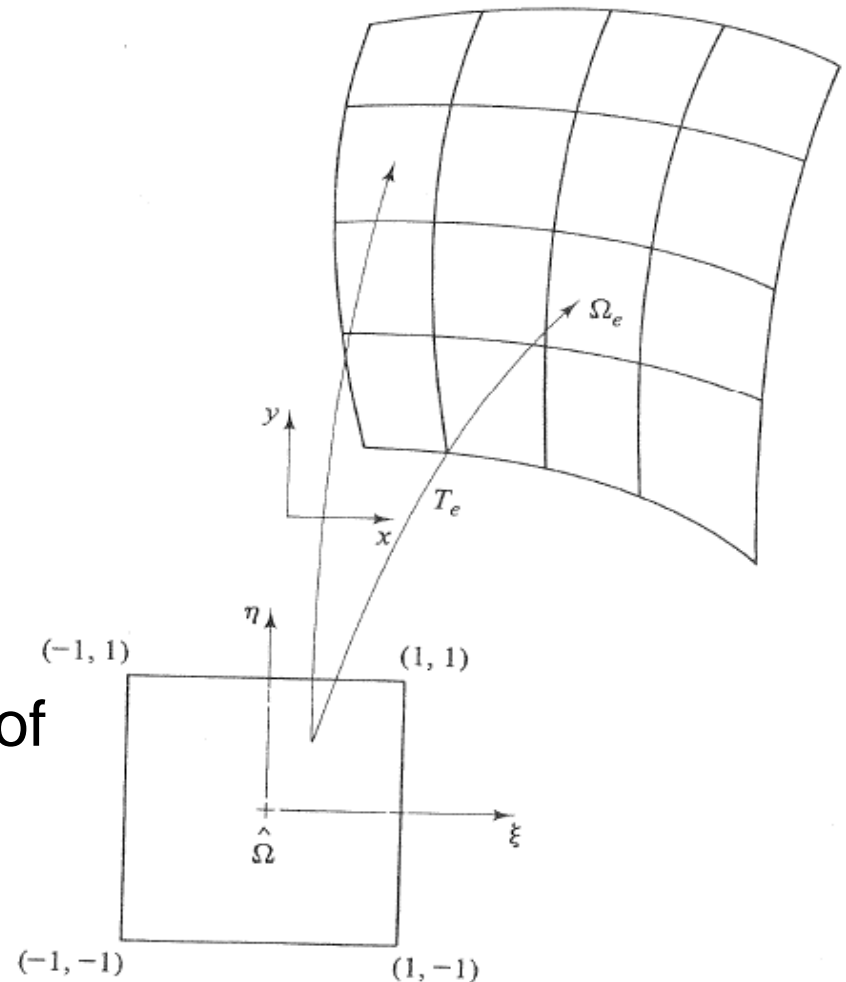
$$dy = \frac{\partial y}{\partial \xi} d\xi + \frac{\partial y}{\partial \eta} d\eta$$

- We can write this as a system of

Eqs:

$$\begin{Bmatrix} dx \\ dy \end{Bmatrix} = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \xi} & \frac{\partial y}{\partial \eta} \end{bmatrix} \begin{Bmatrix} d\xi \\ d\eta \end{Bmatrix}$$

Jacobian matrix $[J]^T$



Transformation equations

- The Jacobian maps segments

$d\xi, d\eta$ from $\hat{\Omega}$ to Ω_e .

$$\begin{Bmatrix} dx \\ dy \end{Bmatrix} = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \xi} & \frac{\partial y}{\partial \eta} \end{bmatrix} \begin{Bmatrix} d\xi \\ d\eta \end{Bmatrix}$$

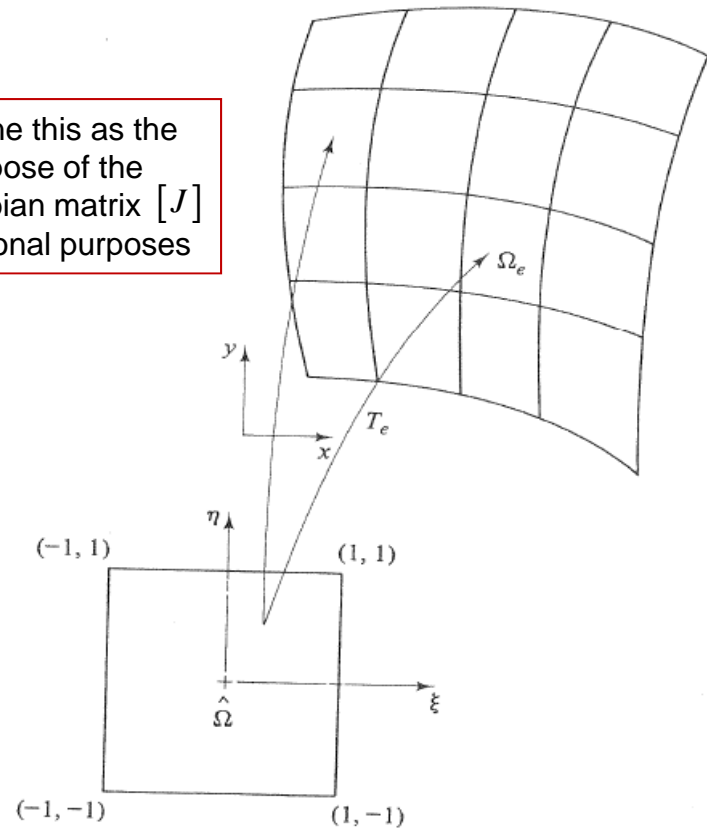
Jacobian matrix $[J]^T$

We define this as the transpose of the Jacobian matrix $[J]$ for notational purposes

- Assume the inverse transformation exists:

$$T_e^{-1} : \Omega_e \rightarrow \hat{\Omega}, e = 1, 2, \dots, E$$

$$\begin{Bmatrix} d\xi \\ d\eta \end{Bmatrix} = \frac{1}{|J|} \begin{bmatrix} \frac{\partial y}{\partial \eta} & -\frac{\partial x}{\partial \eta} \\ -\frac{\partial y}{\partial \xi} & \frac{\partial x}{\partial \xi} \end{bmatrix} \begin{Bmatrix} dx \\ dy \end{Bmatrix}$$



Determinant of the Jacobian matrix

$$|J| = \det[J] = \det[J]^T = \frac{\partial x}{\partial \xi} \frac{\partial y}{\partial \eta} - \frac{\partial x}{\partial \eta} \frac{\partial y}{\partial \xi}$$

Transformation equations

$$\xi = \xi(x, y)$$

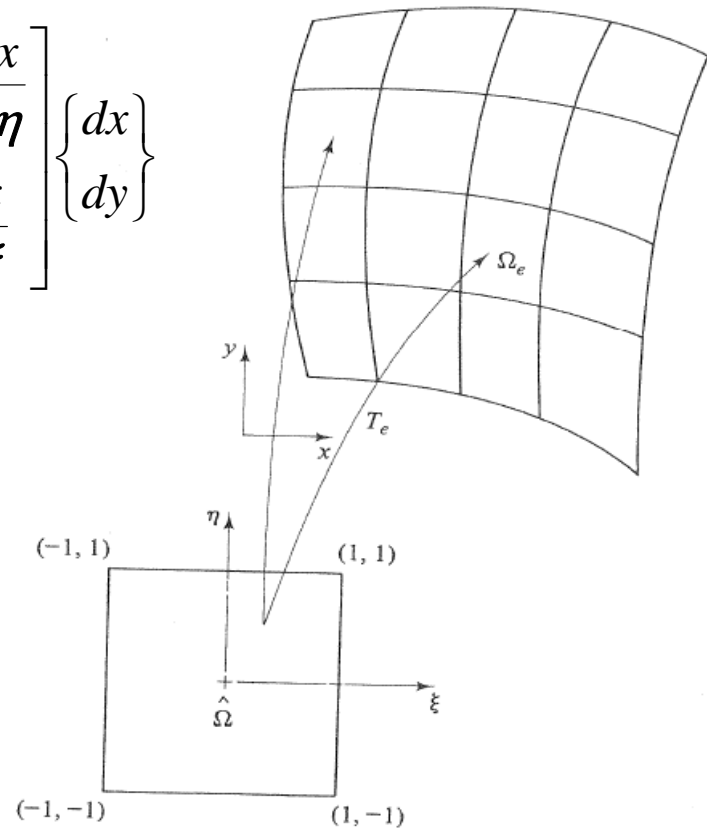
$$T_e^{-1} : \begin{cases} d\xi \\ d\eta \end{cases} = \frac{1}{|J|} \begin{bmatrix} \frac{\partial y}{\partial \eta} & -\frac{\partial x}{\partial \eta} \\ -\frac{\partial y}{\partial \xi} & \frac{\partial x}{\partial \xi} \end{bmatrix} \begin{cases} dx \\ dy \end{cases}$$

$$\eta = \eta(x, y)$$

- From the equations above we can conclude that:

$$\frac{\partial \xi}{\partial x} = \frac{1}{|J|} \frac{\partial y}{\partial \eta}, \quad \frac{\partial \xi}{\partial y} = -\frac{1}{|J|} \frac{\partial x}{\partial \eta},$$

$$\frac{\partial \eta}{\partial x} = -\frac{1}{|J|} \frac{\partial y}{\partial \xi}, \quad \frac{\partial \eta}{\partial y} = \frac{1}{|J|} \frac{\partial x}{\partial \xi}$$



Constructing the transformation T_e

- We are interested for differentiable mappings $\{T_1, T_2, \dots, T_E\}$ of the form $x = x(\xi, \eta)$, $y = y(\xi, \eta)$ for each element that do not create gaps or overlapping between elements and which are easy to construct from the geometry of each element Ω_e .
- Based on an earlier lecture on construction of element basis functions, suppose that we have in place element basis functions $\hat{N}_j(\xi, \eta)$ in $\hat{\Omega}$ such that any function $\hat{g}(\xi, \eta)$ can be approximated as:

$$\hat{g}(\xi, \eta) = \sum_{j=1}^M g_j \hat{N}_j(\xi, \eta), \quad g_j = \hat{g}\left(\underbrace{\xi_j, \eta_j}_{\substack{\text{coordinates of node } j \\ \text{in the master element}}}\right)$$

where M is the number of nodes in the master element.

Constructing the transformation T_e

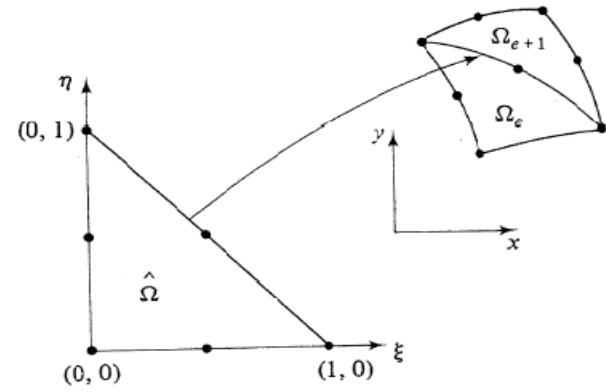
- The finite element basis functions can be used to construct the mapping T_e as follows:

$$x = \sum_{j=1}^M x_j \hat{N}_j(\xi, \eta),$$

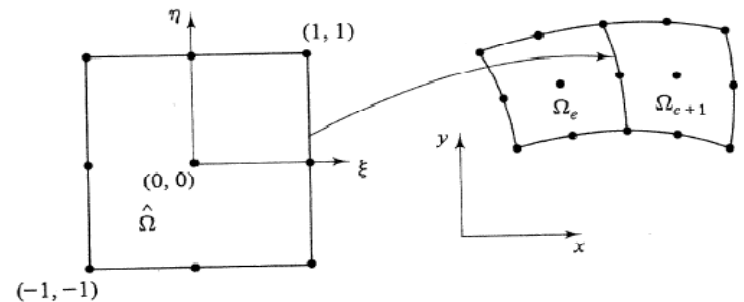
$$y = \sum_{j=1}^M y_j \hat{N}_j(\xi, \eta),$$

where (x_j, y_j) are the coordinates of node j in the element Ω_e .

- Can you convince yourselves that **this mapping does not create gaps or element overlapping?**



Quadratic shape functions



Biquadratic shape functions

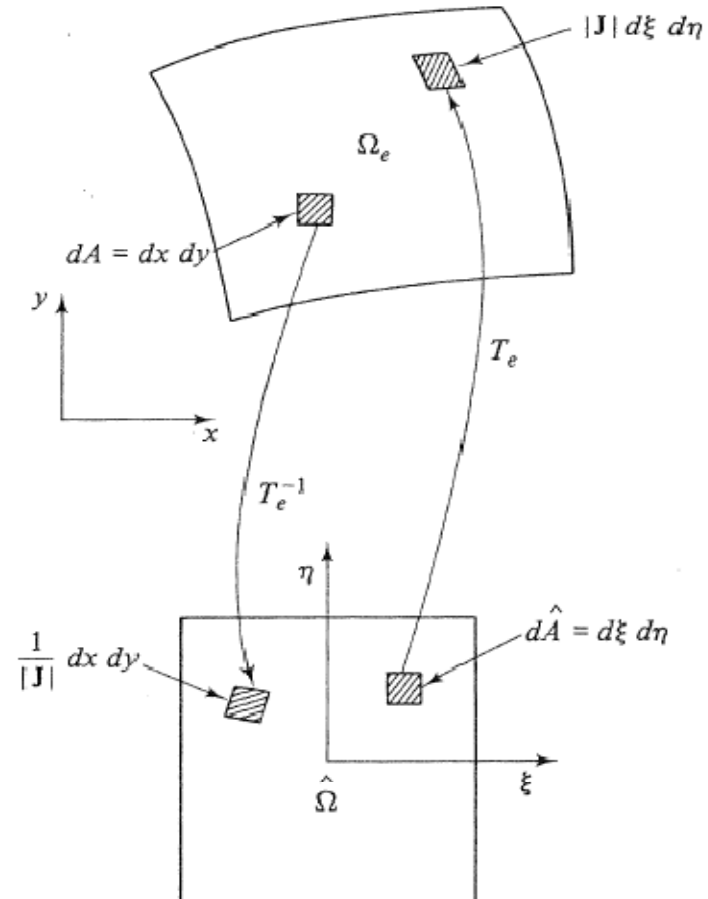
Note that the common side between elements e and $e+1$ is uniquely defined in terms of the nodal coordinates that are shared by the 2 elements.

Geometric interpretation of the Jacobian

- The value of $|J|$ is the ratio of areas of elements at points $x = x(\xi, \eta)$, $y = y(\xi, \eta)$ and (ξ, η) .

$$dA = |J| d\hat{A}$$

- We need to select the element basis functions \hat{N}_i and (x_i, y_i) such that $|J| > 0$ (area cannot be transformed to a line or a point or turned inside out!).
- This guarantees that the mapping T_e is invertible.



Transformation equations

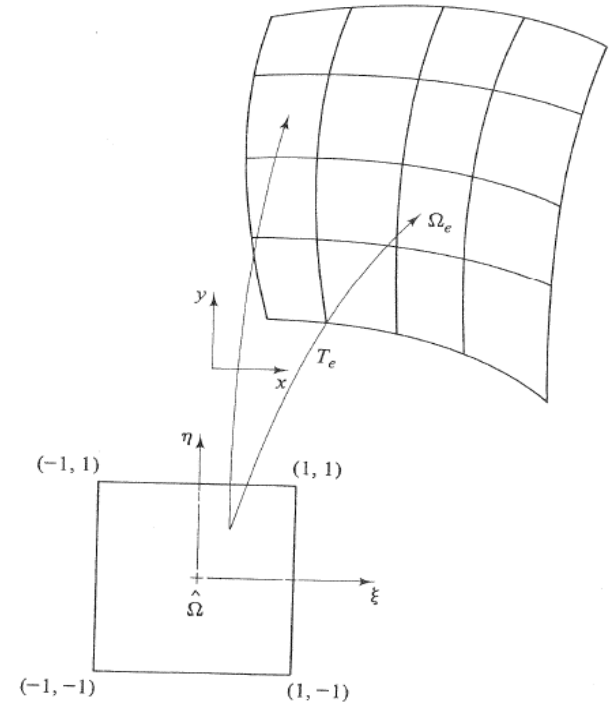
- Let us return to the definition of $|J|$ and the earlier stated transformation equations:

$$|J| = \det J = \frac{\partial x}{\partial \xi} \frac{\partial y}{\partial \eta} - \frac{\partial x}{\partial \eta} \frac{\partial y}{\partial \xi}$$

$$\frac{\partial \xi}{\partial x} = \frac{1}{|J|} \frac{\partial y}{\partial \eta}, \quad \frac{\partial \xi}{\partial y} = -\frac{1}{|J|} \frac{\partial x}{\partial \eta}, \quad \frac{\partial \eta}{\partial x} = -\frac{1}{|J|} \frac{\partial y}{\partial \xi}, \quad \frac{\partial \eta}{\partial y} = \frac{1}{|J|} \frac{\partial x}{\partial \xi}$$

- Applying the mapping:

$$x = \sum_{j=1}^M x_j \widehat{N}_j(\xi, \eta), \quad y = \sum_{j=1}^M y_j \widehat{N}_j(\xi, \eta),$$



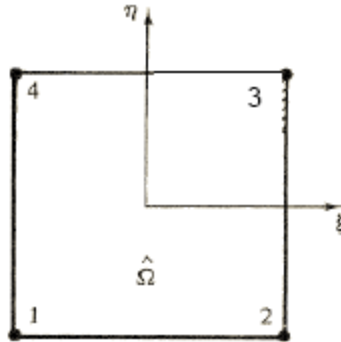
leads to:

$$|J| = \det J = \begin{pmatrix} \sum_{j=1}^M x_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \xi} \\ \sum_{j=1}^M y_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \xi} \end{pmatrix} \begin{pmatrix} \sum_{j=1}^M x_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \eta} \\ \sum_{j=1}^M y_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \eta} \end{pmatrix} - \begin{pmatrix} \sum_{j=1}^M x_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \eta} \\ \sum_{j=1}^M y_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \eta} \end{pmatrix} \begin{pmatrix} \sum_{j=1}^M x_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \xi} \\ \sum_{j=1}^M y_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \xi} \end{pmatrix}$$

$$\frac{\partial \xi}{\partial x} = \frac{1}{|J|} \sum_{j=1}^M y_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \eta}, \quad \frac{\partial \xi}{\partial y} = -\frac{1}{|J|} \sum_{j=1}^M x_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \eta}, \quad \frac{\partial \eta}{\partial x} = -\frac{1}{|J|} \sum_{j=1}^M y_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \xi}, \quad \frac{\partial \eta}{\partial y} = \frac{1}{|J|} \sum_{j=1}^M x_j \frac{\partial \widehat{N}_j(\xi, \eta)}{\partial \xi}$$

Example problem

- Consider the following master element.



- The basis functions for this element are:

$$\hat{N}_1(\xi, \eta) = \frac{1}{4}(1 - \xi)(1 - \eta)$$

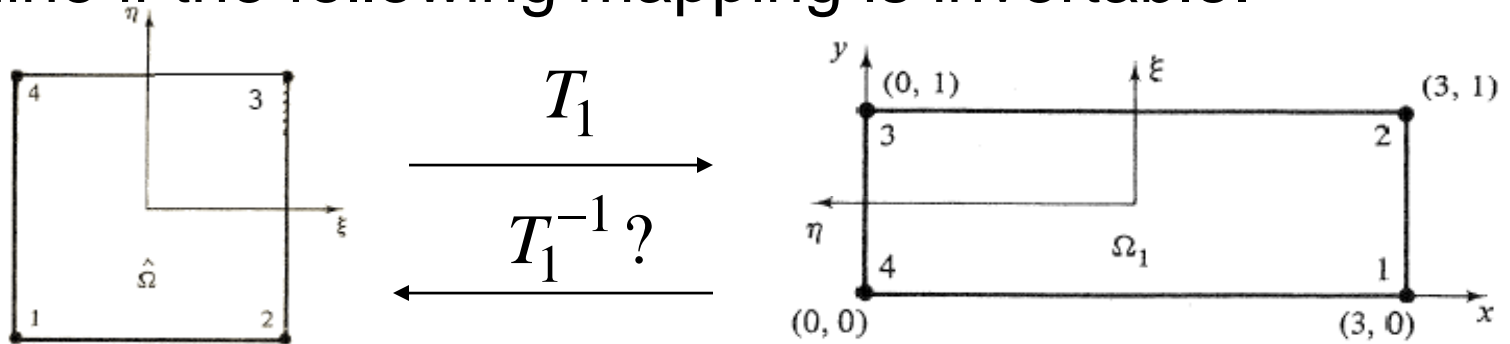
$$\hat{N}_2(\xi, \eta) = \frac{1}{4}(1 + \xi)(1 - \eta)$$

$$\hat{N}_3(\xi, \eta) = \frac{1}{4}(1 + \xi)(1 + \eta)$$

$$\hat{N}_4(\xi, \eta) = \frac{1}{4}(1 - \xi)(1 + \eta)$$

Example problem

- Examine if the following mapping is invertible.



$$x = \sum_{j=1}^M x_j \hat{N}_j(\xi, \eta) = 3\hat{N}_1(\xi, \eta) + 3\hat{N}_2(\xi, \eta) = \frac{3}{2}(1 - \eta)$$

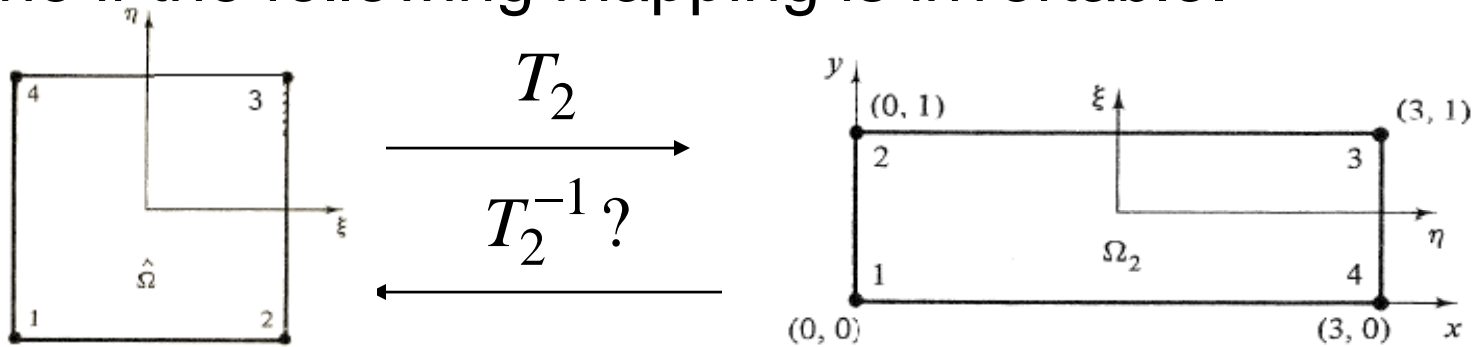
T_1 :

$$y = \sum_{j=1}^M y_j \hat{N}_j(\xi, \eta) = \hat{N}_2(\xi, \eta) + \hat{N}_3(\xi, \eta) = \frac{1}{2}(1 + \xi)$$

- Note that $|J| = \det \begin{bmatrix} 0 & -\frac{3}{2} \\ \frac{1}{2} & 0 \end{bmatrix} = \frac{3}{4} > 0 \Rightarrow$ the map T_1 is invertible.
- $|J|$ is the ratio of the area of Ω_1 to area of $\hat{\Omega}$.

Example problem

- Examine if the following mapping is invertible.



$$x = \sum_{j=1}^M x_j \hat{N}_j(\xi, \eta) = 3\hat{N}_3(\xi, \eta) + 3\hat{N}_4(\xi, \eta) = \frac{3}{2}(1 + \eta)$$

T_2 :

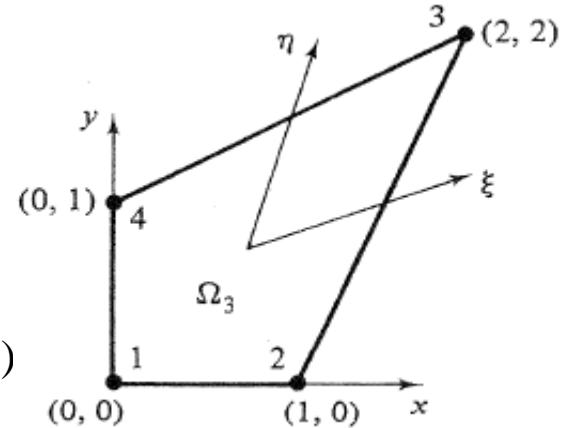
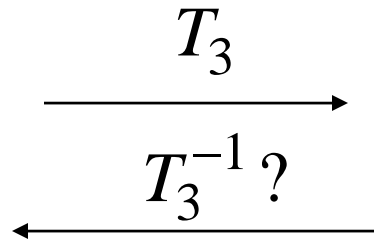
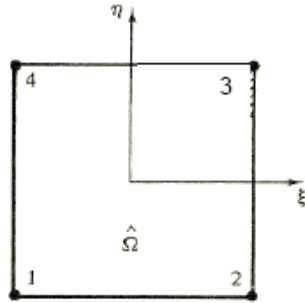
$$y = \sum_{j=1}^M y_j \hat{N}_j(\xi, \eta) = \hat{N}_2(\xi, \eta) + \hat{N}_3(\xi, \eta) = \frac{1}{2}(1 + \xi)$$

- Note that $|J| = \det \begin{bmatrix} 0 & \frac{3}{2} \\ \frac{1}{2} & 0 \end{bmatrix} = -\frac{3}{4} < 0 \Rightarrow$ the map T_2 is not invertible.

- The clockwise numbering of nodes should be avoided.

Example problem

- Examine if the following mapping is invertible.



$$x = \sum_{j=1}^M x_j \widehat{N}_j(\xi, \eta) = \widehat{N}_2(\xi, \eta) + 2\widehat{N}_3(\xi, \eta) = \frac{1}{4}(3 + 3\xi + \eta + \xi\eta)$$

T_3 :

$$y = \sum_{j=1}^M y_j \widehat{N}_j(\xi, \eta) = 2\widehat{N}_3(\xi, \eta) + \widehat{N}_4(\xi, \eta) = \frac{1}{4}(3 + \xi + 3\eta + \xi\eta)$$

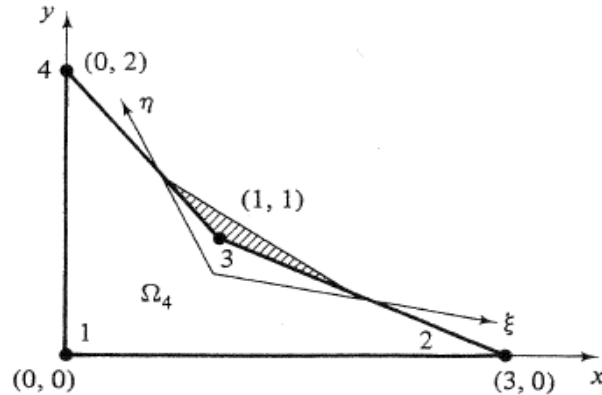
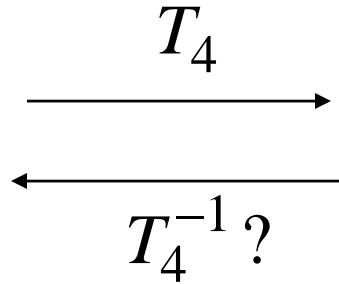
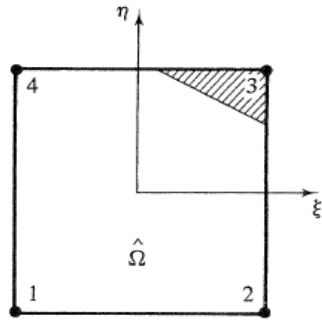
$$|J| = \det \begin{bmatrix} \frac{1}{4}(3+\eta) & \frac{1}{4}(1+\xi) \\ \frac{1}{4}(1+\eta) & \frac{1}{4}(3+\xi) \end{bmatrix} = \frac{1}{2} + \frac{1}{8}\xi + \frac{1}{8}\eta \Rightarrow$$

$|J|$ cannot be zero inside the element so the map T_3 is invertible.

- $|J|$ is smaller near node 1, largest near node 3.

Example problem

- Examine if the following mapping is invertible.



$$x = \sum_{j=1}^M x_j \hat{N}_j(\xi, \eta) = 3\hat{N}_2(\xi, \eta) + \hat{N}_3(\xi, \eta) = \frac{1}{4}(4 + 4\xi - 2\eta - 2\xi\eta)$$

T_4 :

$$y = \sum_{j=1}^M y_j \hat{N}_j(\xi, \eta) = \hat{N}_3(\xi, \eta) + 2\hat{N}_4(\xi, \eta) = \frac{1}{4}(3 - \xi + 3\eta - \xi\eta)$$

$$|J| = \det \begin{bmatrix} \frac{1}{4}(4 - 2\eta) & -\frac{1}{4}(2 + 2\xi) \\ -\frac{1}{4}(1 + \eta) & \frac{1}{4}(3 - \xi) \end{bmatrix} = \frac{1}{8}(5 - 3\xi - 4\eta) \Rightarrow$$

$|J|$ is not > 0 everywhere inside the element so the map T_3 is not invertible. The marked region near node 3 -- above the line $\xi = \frac{5}{3} - \frac{4}{3}\eta$ -- is mapped outside the element.

- All angles of quadrilateral elements need to be $< \pi$.

Finite element equations

- To solve our boundary value problem of interest examined earlier, we need to compute the following matrices and vectors ($1 \leq e \leq E, 1 \leq i, j \leq N_e$):

$$k_{ij}^e = \int_{\Omega_e} \left[k \left(\frac{\partial N_i^e}{\partial x} \frac{\partial N_j^e}{\partial x} + \frac{\partial N_i^e}{\partial y} \frac{\partial N_j^e}{\partial y} \right) + b N_i^e N_j^e \right] dx dy \quad f_i^e = \int_{\Omega_e} f N_i^e dx dy$$

$$P_{ij}^e = \int_{\partial\Omega_{2h}^e} p N_i^e N_j^e ds \quad \gamma_i^e = \int_{\partial\Omega_{2h}^e} \gamma N_i^e ds$$

- k and f are assumed to be functions of (x, y) , whereas p and γ are assumed functions of s defining $\partial\Omega_{2h}(x(s), y(s))$
- It remains to discuss how to compute these using the master element and the transformation T_e .

Master element calculations: Step 1

- We construct the mappings $T_e : \hat{\Omega} \rightarrow \Omega_e, e = 1, 2, \dots, E$ as follows:

$$x = \sum_{j=1}^{N_e} x_j \hat{N}_j(\xi, \eta)$$

$$y = \sum_{j=1}^{N_e} y_j \hat{N}_j(\xi, \eta),$$

- This mapping $T_e : \hat{\Omega} \rightarrow \Omega_e$, is completely defined by the nodal coordinates (x_j, y_j) of element e .

Master element calculations: Step 2

- Let $g(x,y)$ be any function of (x,y) in Ω_e . We can convert g to a function of (ξ,η) defined in $\hat{\Omega}$ as follows:

$$g(x, y) = g(x(\xi, \eta), y(\xi, \eta)) \equiv \hat{g}(\xi, \eta)$$

- The transformation $x(\xi,\eta), y(\xi,\eta)$ is defined in Step 1.
- The element shape functions $N_j^e(x, y)$ are simply obtained from $\hat{N}_j(\xi,\eta)$ as follows:

$$N_j^e(x, y) = \hat{N}_j(x(\xi, \eta), y(\xi, \eta)), \quad j = 1, 2, \dots, N_e$$

Master element calculations: Step 2

- The derivatives of $N_j^e(x, y)$ with respect to x and y are:

$$\frac{\partial N_j^e(x, y)}{\partial x} = \frac{\partial \hat{N}_j}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_j}{\partial \eta} \frac{\partial \eta}{\partial x}, \quad \frac{\partial x}{\partial \xi} = \sum_{k=1}^{N_i} x_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \xi}, \quad \frac{\partial x}{\partial \eta} = \sum_{k=1}^{N_i} x_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \eta}$$

$$\frac{\partial N_j^e(x, y)}{\partial y} = \frac{\partial \hat{N}_j}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_j}{\partial \eta} \frac{\partial \eta}{\partial y} \quad \text{where:} \quad \frac{\partial y}{\partial \eta} = \sum_{k=1}^{N_i} y_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \eta}, \quad \frac{\partial y}{\partial \xi} = \sum_{k=1}^{N_i} y_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \xi}$$

- Finally, with $|J|$ the Jacobian of T_e , using

$$\frac{\partial \xi}{\partial x} = \frac{1}{|J|} \sum_{k=1}^M y_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \eta}, \quad \frac{\partial \xi}{\partial y} = -\frac{1}{|J|} \sum_{k=1}^M x_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \eta}, \quad \frac{\partial \eta}{\partial x} = -\frac{1}{|J|} \sum_{k=1}^M y_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \xi}, \quad \frac{\partial \eta}{\partial y} = \frac{1}{|J|} \sum_{k=1}^M x_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \xi}$$

we conclude:

$$\frac{\partial N_j^e(x, y)}{\partial x} = \frac{1}{|J|} \left\{ \frac{\partial \hat{N}_j}{\partial \xi} \sum_{k=1}^M y_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \eta} - \frac{\partial \hat{N}_j}{\partial \eta} \sum_{k=1}^M y_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \xi} \right\},$$

$$\frac{\partial N_j^e(x, y)}{\partial y} = \frac{1}{|J|} \left\{ -\frac{\partial \hat{N}_j}{\partial \xi} \sum_{k=1}^M x_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \eta} + \frac{\partial \hat{N}_j}{\partial \eta} \sum_{k=1}^M x_k \frac{\partial \hat{N}_k(\xi, \eta)}{\partial \xi} \right\}$$

Note regarding convergence rates

- We had presented earlier that the asymptotic rates of convergence are only valid if $w_h^e(x, y) = \sum_{j=1}^{N_e} w_j N_j^e(x, y)$

contains complete polynomials of degree k for smooth interpolated functions.

- However, with the transformation

$$N_j^e(x, y) = \widehat{N}_j(x(\xi, \eta), y(\xi, \eta)), \quad j = 1, 2, \dots, N_e$$

the shape functions $N_j^e(x, y)$ may not even be polynomials!

- It can be shown that if the master basis functions $\widehat{N}_j(\xi, \eta)$ contain complete polynomials of degree k and $|J| > 0$, the error estimates still hold.

Master element calculations: Step 3

- The domain integrals of interest now take the form:

$$\int_{\Omega} g(x, y) dx dy = \int_{\hat{\Omega}} g(x(\xi, \eta), y(\xi, \eta)) |J| d\xi d\eta \equiv \int_{\hat{\Omega}} \underbrace{\hat{g}(\xi, \eta)}_{\hat{G}(\xi, \eta)} |J| d\xi d\eta$$

- Our integrals are now in the domain of $\hat{\Omega}$

$$\int_{\Omega} g(x, y) dx dy = \int_{\hat{\Omega}} g(x(\xi, \eta), y(\xi, \eta)) |J| d\xi d\eta \equiv \int_{\hat{\Omega}} \hat{G}(\xi, \eta) d\xi d\eta$$

- As in 1D, we use Gauss integration (to be discussed):

$$\int_{\hat{\Omega}} \hat{G}(\xi, \eta) d\xi d\eta \approx \sum_{i=1}^{N_i} \hat{G}(\underbrace{\xi_i, \eta_i}_{\text{Gauss integration points}}) \underbrace{w_i}_{\text{Gauss integration weights}}$$

Master element calculations: Step 4

- In numerical calculation of e.g. the element stiffness,

$$k_{ij}^e = \int_{\Omega_e} \left[k \left(\frac{\partial N_i^e}{\partial x} \frac{\partial N_j^e}{\partial x} + \frac{\partial N_i^e}{\partial y} \frac{\partial N_j^e}{\partial y} \right) + b N_i^e N_j^e \right] dx dy$$

we often need to consider functions $k(x,y)$. It is common to use the finite element interpolant to approximate this function, i.e.

$$k_h(x, y) = \sum_{j=1}^{N_e} \underbrace{k_j}_{k(x_j, y_j)} \quad N_j^e(x, y) = \sum_{j=1}^{N_e} k_j \widehat{N}_j(\xi, \eta)$$

This requires only the nodal values of k . Similar calculation can be done for integrals involving $b(x,y)$ and $f(x,y)$.

Master element calculations: Step 5

- We need to consider a number of integrals that involve boundary integration:

$$P_{ij}^e = \int_{\partial\Omega_{2h}^e} p N_i^e N_j^e ds \qquad \gamma_i^e = \int_{\partial\Omega_{2h}^e} \gamma N_i^e ds$$

- Lets introduce the restriction of the basis functions $\hat{N}_j(\xi, \eta)$ on the side $\xi = 1$. We denote them as: $\hat{\theta}(\eta) \equiv \hat{N}_j(1, \eta)$, $j = 1, 2, \dots, N_e$.

These functions are only non-zero for nodes on the boundary $\xi = 1$. The 2 integrals above can be computed as:

$$P_{ij}^e = \int_{\partial\Omega_{2h}^e} p N_i^e N_j^e ds = \int_{-1}^1 \hat{p}(\eta) \hat{\theta}_i(\eta) \hat{\theta}_j(\eta) |j(\eta)| d\eta$$

$$\gamma_i^e = \int_{\partial\Omega_{2h}^e} \gamma N_i^e ds = \int_{-1}^1 \hat{\gamma}(\eta) \hat{\theta}_i(\eta) |j(\eta)| d\eta$$

$|j(\eta)|$ is the Jacobian of the transformation of η to s
 $ds = |j(\eta)| d\eta \quad \frac{ds}{dx} = \left[dx^2 + dy^2 \right]^{1/2}$
 $|j(\eta)| = \left[\left(\frac{\partial x(1, \eta)}{\partial \eta} \right)^2 + \left(\frac{\partial y(1, \eta)}{\partial \eta} \right)^2 \right]^{1/2}$

Notes: (a) Use FE interpolants of $\hat{p}(\eta), \hat{\gamma}(\eta)$ (b) Use 1D Gauss integration.

Example problem using 3-node triangular elements

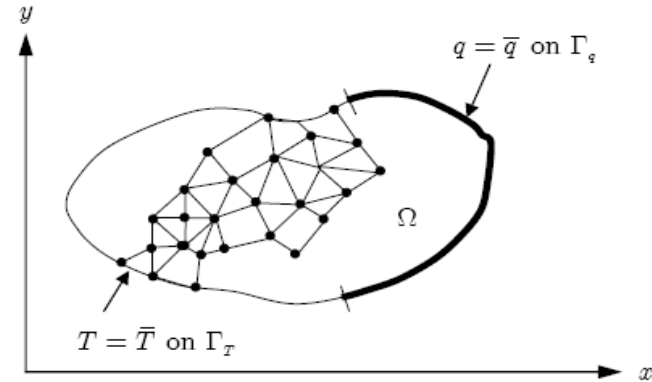
- Consider 2D heat conduction using the general problem discussed in [an earlier lecture](#).

- We are computing $T(x,y)$ such that:

$$-\nabla \cdot (D\nabla T) = f(x, y)$$

$$T = \bar{T} \text{ on } \Gamma_T$$

$$q = -D\nabla T \cdot n = \bar{q} \text{ on } \Gamma_q$$



- Recall that the final weak form is:

$$\{w\}^T \sum_{e=1}^{nel} \underbrace{[L^e]^T \int_{\Omega^e} [B^e]^T [D^e] [B^e] d\Omega [L^e]}_{K^e} \{d\} = \{w\}^T \sum_{e=1}^{nel} [L^e]^T \left(\underbrace{\int_{\Omega^e} [N^e]^T f d\Omega}_{\{f_{\Omega}^e\}} - \underbrace{\int_{\Gamma_q^e} [N^e]^T \bar{q} d\Gamma}_{\{f_{\Gamma}^e\}} \right), \forall \{w_F\}$$

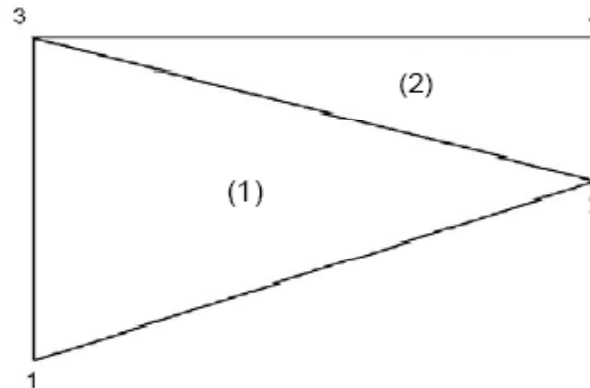
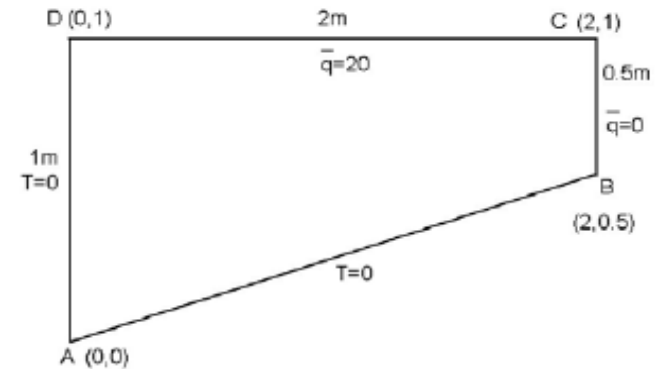
$$[B^e] = \begin{bmatrix} \frac{\partial N_1^e}{\partial x} & \frac{\partial N_2^e}{\partial x} & \dots & \frac{\partial N_{nen}^e}{\partial x} \\ \frac{\partial N_1^e}{\partial y} & \frac{\partial N_2^e}{\partial y} & \dots & \frac{\partial N_{nen}^e}{\partial y} \end{bmatrix}$$

$$[N^e] = [N_1^e \quad N_2^e \quad \dots \quad N_{nen}^e]$$

$$\begin{bmatrix} K_E & K_{EF} \\ K_{FE} & K_F \end{bmatrix} \begin{bmatrix} \bar{d}_E \\ d_F \end{bmatrix} = \begin{bmatrix} f_E + r_E \\ f_F \end{bmatrix}$$

Example problem using 3-node triangular elements

- Consider the heat conduction problem in the figure. We assume isotropic conductivity $k=5 \text{ W}^\circ\text{C}^{-1}$ and a source term $f=6 \text{ Wm}^{-2}$. The BCs are as shown.
- We want to compute the temperature field using 2 linear triangular elements.



Example problem : Stiffness, element 1

- Using expressions derived before for triangular elements, we can write:

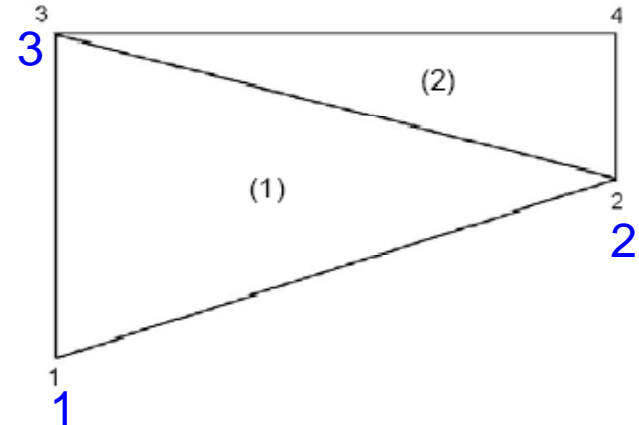
$$\underbrace{[B^e]}_{\text{constant matrix}} = \begin{bmatrix} \frac{\partial N_1^e}{\partial x} & \frac{\partial N_2^e}{\partial x} & \frac{\partial N_3^e}{\partial x} \\ \frac{\partial N_1^e}{\partial y} & \frac{\partial N_2^e}{\partial y} & \frac{\partial N_3^e}{\partial y} \end{bmatrix} = \frac{1}{2A^e} \begin{bmatrix} y_2^e - y_3^e & y_3^e - y_1^e & y_1^e - y_2^e \\ x_3^e - x_2^e & x_1^e - x_3^e & x_2^e - x_1^e \end{bmatrix}$$

$$2A^e = (x_2^e y_3^e - x_3^e y_2^e) - (x_1^e y_3^e - x_3^e y_1^e) + (x_1^e y_2^e - x_2^e y_1^e)$$

$$[K^e] = \int_{\Omega^e} [B^e]^T [D^e] [B^e] d\Omega = [B^e]^T k [B^e] A^e$$

- For $e=1$: $[B^1] = \frac{1}{2} \begin{bmatrix} -0.5 & 1 & -0.5 \\ -2 & 0 & 2 \end{bmatrix}$

$$K^1 = B^{1T} B^1 k A^1 = \begin{bmatrix} 5.3125 & -0.625 & -4.6875 \\ -0.625 & 1.25 & -0.625 \\ -4.6875 & -0.625 & 5.3125 \end{bmatrix} \begin{matrix} 1 \\ 2 \\ 3 \end{matrix}$$



Example problem : Stiffness, element 2

- Using expressions derived before for triangular elements, we can write:

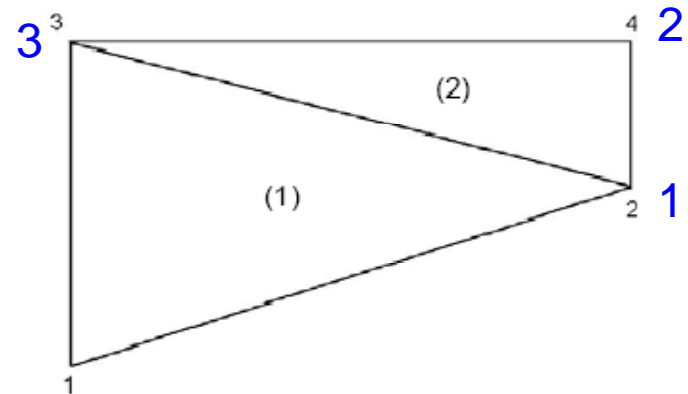
$$\underbrace{B^e}_{\text{constant matrix}} = \begin{bmatrix} \frac{\partial N_1^e}{\partial x} & \frac{\partial N_2^e}{\partial x} & \frac{\partial N_3^e}{\partial x} \\ \frac{\partial N_1^e}{\partial y} & \frac{\partial N_2^e}{\partial y} & \frac{\partial N_3^e}{\partial y} \end{bmatrix} = \frac{1}{2A^e} \begin{bmatrix} y_2^e - y_3^e & y_3^e - y_1^e & y_1^e - y_2^e \\ x_3^e - x_2^e & x_1^e - x_3^e & x_2^e - x_1^e \end{bmatrix}$$

$$2A^e = (x_2^e y_3^e - x_3^e y_2^e) - (x_1^e y_3^e - x_3^e y_1^e) + (x_1^e y_2^e - x_2^e y_1^e)$$

$$K^e = \int_{\Omega^e} B^{eT} [D^e] B^e d\Omega = B^{eT} B k A^e$$

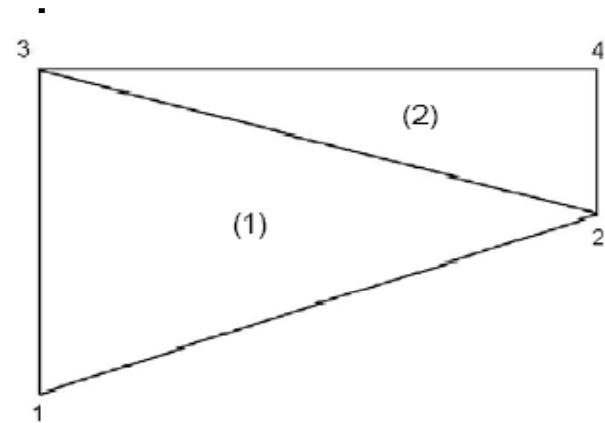
- For $e=2$: $B^2 = \begin{bmatrix} 0 & 0.5 & -0.5 \\ -2 & 2 & 0 \end{bmatrix}$

$$[K^2] = [B^2]^T [B^2] k A^2 = \begin{bmatrix} 10 & -10 & 0 \\ -10 & 10.625 & -0.625 \\ 0 & -0.625 & 0.625 \end{bmatrix} \begin{matrix} 2 \\ 4 \\ 3 \end{matrix}$$



Example problem : Assembled stiffness

- Assembling the 2 stiffness matrices



$$[K] = [L^1]^T [K^1] [L^1] + [L^2]^T [K^2] [L^2] = \begin{bmatrix} 5.3125 & -0.625 & -4.6875 & 0 \\ -0.625 & 11.25 & -0.625 & -10 \\ -4.6875 & -0.625 & 5.9375 & -0.625 \\ 0 & -10 & -0.625 & 10.625 \end{bmatrix} \begin{matrix} 1 \\ 2 \\ 3 \\ 4 \end{matrix}$$

Example problem : Load calculations – e=1,2

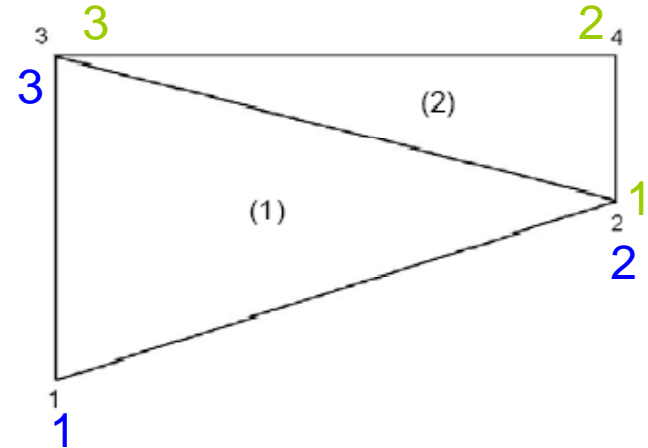
- The shape functions for 3-node triangular elements are

$$[N^e] = [N_1^e \quad N_2^e \quad N_3^e]$$

$$N_1^e = \frac{1}{2A^e} (x_2^e y_3^e - x_3^e y_2^e + (y_2^e - y_3^e)x + (x_3^e - x_2^e)y)$$

$$N_2^e = \frac{1}{2A^e} (x_3^e y_1^e - x_1^e y_3^e + (y_3^e - y_1^e)x + (x_1^e - x_3^e)y)$$

$$N_3^e = \frac{1}{2A^e} (x_1^e y_2^e - x_2^e y_1^e + (y_1^e - y_2^e)x + (x_2^e - x_1^e)y)$$



$$2A^e = (x_2^e y_3^e - x_3^e y_2^e) - (x_1^e y_3^e - x_3^e y_1^e) + (x_1^e y_2^e - x_2^e y_1^e)$$

$$\{f_\Omega^e\} = f \int_{\Omega^e} \{N^e\}^T d\Omega = \frac{fA^e}{3} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

This expression can be proved for constant f!

- For e=1:

$$\{f^1\} = \frac{fA^1}{3} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} = \frac{6 \times 1}{3} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} = \begin{bmatrix} 2 \\ 2 \\ 2 \end{bmatrix}$$

- For e=2:

$$\{f^2\} = \frac{fA^2}{3} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} = \frac{6 \times 0.5}{3} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

$$\{f\} = [L^1]^T \{f^1\} + [L^2]^T \{f^2\} = \begin{bmatrix} 2 \\ 2+1 \\ 2+1 \\ 1 \end{bmatrix} = \begin{bmatrix} 2 \\ 3 \\ 3 \\ 1 \end{bmatrix}$$

Example problem: Applying natural BCs

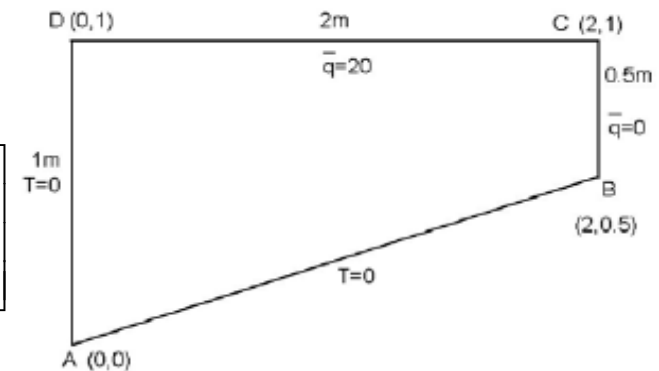
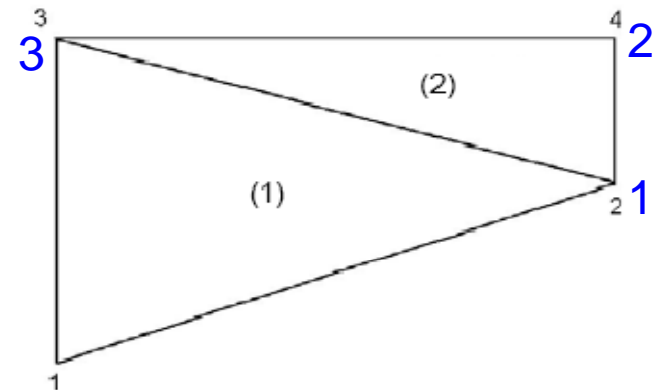
- For elements with segments on Γ_q , we need to compute

$$\{f_\Gamma^e\} = - \int_{\Gamma_q^e} [N^e]^T \bar{q} d\Gamma$$

- We only need to worry about $e=2$

$$[N^2]^T|_{y=1} = \begin{bmatrix} \frac{1}{2A^2}(x_2^2 y_3^2 - x_3^2 y_2^2 + (y_2^2 - y_3^2)x + (x_3^2 - x_2^2)y) \\ \frac{1}{2A^2}(x_3^2 y_1^2 - x_1^2 y_3^2 + (y_3^2 - y_1^2)x + (x_1^2 - x_3^2)y) \\ \frac{1}{2A^2}(x_1^2 y_2^2 - x_2^2 y_1^2 + (y_1^2 - y_2^2)x + (x_2^2 - x_1^2)y) \end{bmatrix}_{y=1} = \begin{bmatrix} 0 \\ 0.5x \\ -0.5x + 1 \end{bmatrix}$$

- Thus: $\{f_\Gamma^e\} = - \int_0^2 \begin{bmatrix} 0 \\ 0.5x \\ -0.5x + 1 \end{bmatrix} 20 dx = \begin{bmatrix} 0 \\ -20 \\ -20 \end{bmatrix}$



Example problem: Final system of equations

- Assembling the load vector gives the following:

$$\{f\} = \{f_\Gamma\} + \{f_\Omega\} + \{r\} = \begin{bmatrix} 2 \\ 3 \\ -17 \\ -19 \end{bmatrix} + \begin{bmatrix} r_1 \\ r_2 \\ r_3 \\ 0 \end{bmatrix} = \begin{bmatrix} r_1 + 2 \\ r_2 + 3 \\ r_3 - 17 \\ -19 \end{bmatrix}$$

r_1, r_2, r_3 , are unknown heat fluxes at the nodes with prescribed temperature (short of reaction forces!)

- The final system of equations is:

$$\begin{array}{ccc|c|c} 5.3125 & -0.625 & -4.6875 & 0 & 0 \\ -0.625 & 11.25 & -0.625 & -10 & 0 \\ -4.6875 & -0.625 & 5.9375 & -0.625 & 0 \\ \hline 0 & -10 & -0.625 & 10.625 & T_4 \end{array} = \begin{bmatrix} r_1 + 2 \\ r_2 + 3 \\ r_3 - 17 \\ -19 \end{bmatrix}$$

- Partitioning this system gives: $T_4 = -\frac{19}{10.625} = -1.788$

- 'Reaction fluxes': $\begin{bmatrix} 5.3125 & -0.625 & -4.6875 & 0 \\ -0.625 & 11.25 & -0.625 & -10 \\ -4.6875 & -0.625 & 5.9375 & -0.625 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \\ -1.788 \end{bmatrix} = \begin{bmatrix} r_1 + 2 \\ r_2 + 3 \\ r_3 - 17 \end{bmatrix} \Rightarrow \begin{bmatrix} r_1 \\ r_2 \\ r_3 \end{bmatrix} = \begin{bmatrix} -2 \\ 14.88 \\ 18.1175 \end{bmatrix}$

Example problem: Postprocessing

- The DOF in each element (nodal temperatures) are now given as:

$$\{d^1\} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}, \{d^2\} = \begin{bmatrix} 0 \\ -1.788 \\ 0 \end{bmatrix}$$

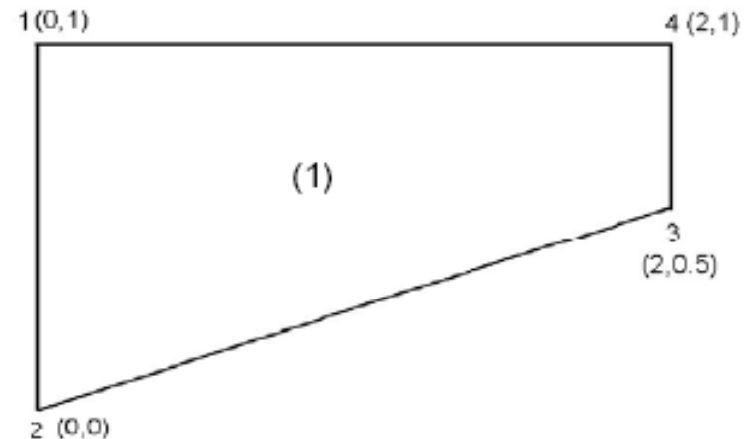
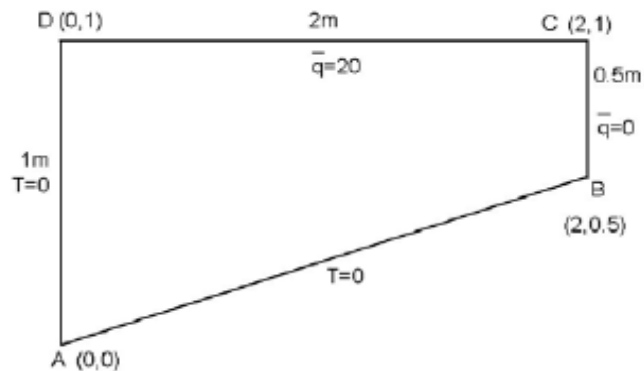
- The flux vectors can then be computed for each element as:

$$\{q^1\} \equiv \begin{bmatrix} q_x^1 \\ q_y^1 \end{bmatrix} = -k[B^1]\{d^1\} = -5 \frac{1}{2} \begin{bmatrix} -0.5 & 1 & -0.5 \\ -2 & 0 & 2 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\{q^2\} \equiv \begin{bmatrix} q_x^2 \\ q_y^2 \end{bmatrix} = -k[B^2]\{d^2\} = -5 \begin{bmatrix} 0 & 0.5 & -0.5 \\ -2 & 2 & 2 \end{bmatrix} \begin{bmatrix} 0 \\ -1.788 \\ 0 \end{bmatrix} = \begin{bmatrix} 4.47 \\ 17.88 \end{bmatrix}$$

Example problem using 4-node quadrilateral elements

- We repeat the calculations of the earlier example using one 4-node quadrilateral element. We use again 'matrix notation'.



- The element coordinates are:

$$\begin{bmatrix} x^e & y^e \end{bmatrix} = \begin{bmatrix} x_1^e & y_1^e \\ x_2^e & y_2^e \\ x_3^e & y_3^e \\ x_4^e & y_4^e \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 0 & 0 \\ 2 & 0.5 \\ 2 & 1 \end{bmatrix}$$

Example problem using 4-node quadrilateral elements

- The basis functions in natural coordinates are:

$$\hat{N}_1(\xi, \eta) = \frac{1}{4}(1-\xi)(1-\eta)$$

$$\hat{N}_2(\xi, \eta) = \frac{1}{4}(1+\xi)(1-\eta)$$

$$\hat{N}_3(\xi, \eta) = \frac{1}{4}(1+\xi)(1+\eta)$$

$$\hat{N}_4(\xi, \eta) = \frac{1}{4}(1-\xi)(1+\eta)$$

- We need to compute the B^e matrix relating derivatives wrt x and y and nodal values.

$$\begin{bmatrix} \frac{\partial T^e}{\partial x} \\ \frac{\partial T^e}{\partial y} \end{bmatrix} = \underbrace{\begin{bmatrix} \frac{\partial \hat{N}_1^e}{\partial x} & \frac{\partial \hat{N}_2^e}{\partial x} & \frac{\partial \hat{N}_3^e}{\partial x} & \frac{\partial \hat{N}_4^e}{\partial x} \\ \frac{\partial \hat{N}_1^e}{\partial y} & \frac{\partial \hat{N}_2^e}{\partial y} & \frac{\partial \hat{N}_3^e}{\partial y} & \frac{\partial \hat{N}_4^e}{\partial y} \end{bmatrix}}_{[B^e]} \begin{bmatrix} T_1^e \\ T_2^e \\ T_3^e \\ T_4^e \end{bmatrix}$$

$$\begin{bmatrix} \frac{\partial T^e}{\partial x} \\ \frac{\partial T^e}{\partial y} \end{bmatrix} = \begin{bmatrix} \frac{\partial \hat{N}_1^e}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_1^e}{\partial \eta} \frac{\partial \eta}{\partial x} & \frac{\partial \hat{N}_2^e}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_2^e}{\partial \eta} \frac{\partial \eta}{\partial x} & \frac{\partial \hat{N}_3^e}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_3^e}{\partial \eta} \frac{\partial \eta}{\partial x} & \frac{\partial \hat{N}_4^e}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_4^e}{\partial \eta} \frac{\partial \eta}{\partial x} \\ \frac{\partial \hat{N}_1^e}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_1^e}{\partial \eta} \frac{\partial \eta}{\partial y} & \frac{\partial \hat{N}_2^e}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_2^e}{\partial \eta} \frac{\partial \eta}{\partial y} & \frac{\partial \hat{N}_3^e}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_3^e}{\partial \eta} \frac{\partial \eta}{\partial y} & \frac{\partial \hat{N}_4^e}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_4^e}{\partial \eta} \frac{\partial \eta}{\partial y} \end{bmatrix} \begin{bmatrix} T_1^e \\ T_2^e \\ T_3^e \\ T_4^e \end{bmatrix}$$

Example problem using 4-node quadrilateral elements

$$[B^e] = \begin{bmatrix} \frac{\partial \hat{N}_1^e}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_1^e}{\partial \eta} \frac{\partial \eta}{\partial x} & \frac{\partial \hat{N}_2^e}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_2^e}{\partial \eta} \frac{\partial \eta}{\partial x} & \frac{\partial \hat{N}_3^e}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_3^e}{\partial \eta} \frac{\partial \eta}{\partial x} & \frac{\partial \hat{N}_4^e}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \hat{N}_4^e}{\partial \eta} \frac{\partial \eta}{\partial x} \\ \frac{\partial \hat{N}_1^e}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_1^e}{\partial \eta} \frac{\partial \eta}{\partial y} & \frac{\partial \hat{N}_2^e}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_2^e}{\partial \eta} \frac{\partial \eta}{\partial y} & \frac{\partial \hat{N}_3^e}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_3^e}{\partial \eta} \frac{\partial \eta}{\partial y} & \frac{\partial \hat{N}_4^e}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \hat{N}_4^e}{\partial \eta} \frac{\partial \eta}{\partial y} \end{bmatrix}$$

- We can write this matrix as:

$$[B^e] = \begin{bmatrix} \frac{\partial \xi}{\partial x} & \frac{\partial \eta}{\partial x} \\ \frac{\partial \xi}{\partial y} & \frac{\partial \eta}{\partial y} \end{bmatrix} \begin{bmatrix} \frac{\partial \hat{N}_1^e}{\partial \xi} & \frac{\partial \hat{N}_2^e}{\partial \xi} & \frac{\partial \hat{N}_3^e}{\partial \xi} & \frac{\partial \hat{N}_4^e}{\partial \xi} \\ \frac{\partial \hat{N}_1^e}{\partial \eta} & \frac{\partial \hat{N}_2^e}{\partial \eta} & \frac{\partial \hat{N}_3^e}{\partial \eta} & \frac{\partial \hat{N}_4^e}{\partial \eta} \end{bmatrix} = \frac{1}{|J^e|} \begin{bmatrix} \frac{\partial y^e}{\partial \eta} & -\frac{\partial y^e}{\partial \xi} \\ -\frac{\partial x^e}{\partial \eta} & \frac{\partial x^e}{\partial \xi} \end{bmatrix} \begin{bmatrix} \frac{\partial \hat{N}_1^e}{\partial \xi} & \frac{\partial \hat{N}_2^e}{\partial \xi} & \frac{\partial \hat{N}_3^e}{\partial \xi} & \frac{\partial \hat{N}_4^e}{\partial \xi} \\ \frac{\partial \hat{N}_1^e}{\partial \eta} & \frac{\partial \hat{N}_2^e}{\partial \eta} & \frac{\partial \hat{N}_3^e}{\partial \eta} & \frac{\partial \hat{N}_4^e}{\partial \eta} \end{bmatrix}$$

This comes from an earlier derivation:

$$\begin{Bmatrix} d\xi \\ d\eta \end{Bmatrix} = \frac{1}{|J^e|} \begin{bmatrix} \frac{\partial y^e}{\partial \eta} & -\frac{\partial x^e}{\partial \eta} \\ -\frac{\partial y^e}{\partial \xi} & \frac{\partial x^e}{\partial \xi} \end{bmatrix} \begin{Bmatrix} dx \\ dy \end{Bmatrix}$$

Example problem using 4-node quadrilateral elements

$$[B^e] = \frac{1}{|J^e|} \underbrace{\begin{bmatrix} \frac{\partial y^e}{\partial \eta} & -\frac{\partial y^e}{\partial \xi} \\ -\frac{\partial x^e}{\partial \eta} & \frac{\partial x^e}{\partial \xi} \end{bmatrix}}_{\text{Inverse of the Jacobian } [J^e]} \begin{bmatrix} \frac{\partial \hat{N}_1^e}{\partial \xi} & \frac{\partial \hat{N}_2^e}{\partial \xi} & \frac{\partial \hat{N}_3^e}{\partial \xi} & \frac{\partial \hat{N}_4^e}{\partial \xi} \\ \frac{\partial \hat{N}_1^e}{\partial \eta} & \frac{\partial \hat{N}_2^e}{\partial \eta} & \frac{\partial \hat{N}_3^e}{\partial \eta} & \frac{\partial \hat{N}_4^e}{\partial \eta} \end{bmatrix}$$

$$x^e = \sum_{j=1}^4 x_j^e \hat{N}_j(\xi, \eta),$$

- The Jacobian matrix can be computed from: $y^e = \sum_{j=1}^4 y_j^e \hat{N}_j(\xi, \eta),$

$$\underbrace{[J^e]}_{\text{Jacobian matrix of element } e} = \begin{bmatrix} \frac{\partial x^e}{\partial \xi} & \frac{\partial y^e}{\partial \xi} \\ \frac{\partial x^e}{\partial \eta} & \frac{\partial y^e}{\partial \eta} \end{bmatrix} = \begin{bmatrix} \frac{\partial \hat{N}_1^e}{\partial \xi} & \frac{\partial \hat{N}_2^e}{\partial \xi} & \frac{\partial \hat{N}_3^e}{\partial \xi} & \frac{\partial \hat{N}_4^e}{\partial \xi} \\ \frac{\partial \hat{N}_1^e}{\partial \eta} & \frac{\partial \hat{N}_2^e}{\partial \eta} & \frac{\partial \hat{N}_3^e}{\partial \eta} & \frac{\partial \hat{N}_4^e}{\partial \eta} \end{bmatrix} \begin{bmatrix} x_1^e & y_1^e \\ x_2^e & y_2^e \\ x_3^e & y_3^e \\ x_4^e & y_4^e \end{bmatrix}$$

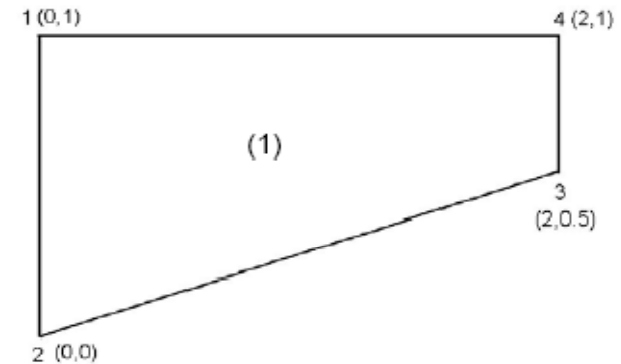
Example problem using 4-node quadrilateral elements

- For element 1:

$$[J^1] = \frac{1}{4} \begin{bmatrix} \eta-1 & 1-\eta & 1+\eta & -\eta-1 \\ \xi-1 & -\xi-1 & 1+\xi & 1-\xi \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 0 & 0 \\ 2 & 0.5 \\ 2 & 1 \end{bmatrix} = \begin{bmatrix} 0 & 0.125\eta - 0.375 \\ 1 & 0.125\xi + 0.125 \end{bmatrix}$$

$$[J^1]^{-1} = \begin{bmatrix} \frac{1+\xi}{3-\eta} & 1 \\ \frac{8}{\eta-3} & 0 \end{bmatrix}, \det J^1 = -0.125\eta + 0.375$$

$$[B^1] = \begin{bmatrix} \frac{1+\xi}{3-\eta} & 1 \\ \frac{8}{\eta-3} & 0 \end{bmatrix} \frac{1}{4} \begin{bmatrix} \eta-1 & 1-\eta & 1+\eta & -\eta-1 \\ \xi-1 & -\xi-1 & 1+\xi & 1-\xi \end{bmatrix}$$

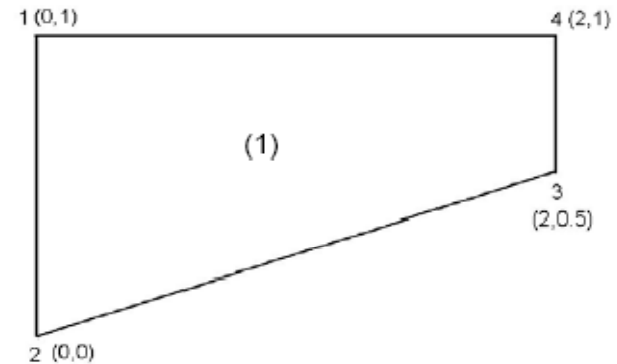


- The element stiffness is: $[K^1] = \int_{\Omega} [B^1]^T [D^1] [B^1] d\Omega = \int_{-1}^1 \int_{-1}^1 [B^1]^T \underbrace{[D^1]}_{k[I]} [B^1] |J^1| d\xi d\eta \Rightarrow$
 $[K^1] = \sum_{i=1}^{N_{in}=2} \sum_{j=1}^{N_{in}=2} k \left([B^1]^T [B^1] |J^1| \right)_{\xi=\xi_i, \eta=\eta_j} W_i W_j, (\xi_i, \eta_j) = \left(\pm \frac{1}{\sqrt{3}}, \pm \frac{1}{\sqrt{3}} \right), (W_i, W_j) = (1, 1)$

Example problem using 4-node quadrilateral elements

- The element stiffness is finally given as:

$$[K^1] = \begin{bmatrix} 4.76 & -3.51 & -2.98 & 1.73 \\ -3.51 & 4.13 & 1.73 & -2.36 \\ -2.98 & 1.73 & 6.54 & -5.29 \\ 1.73 & -2.36 & -5.29 & 5.91 \end{bmatrix}$$

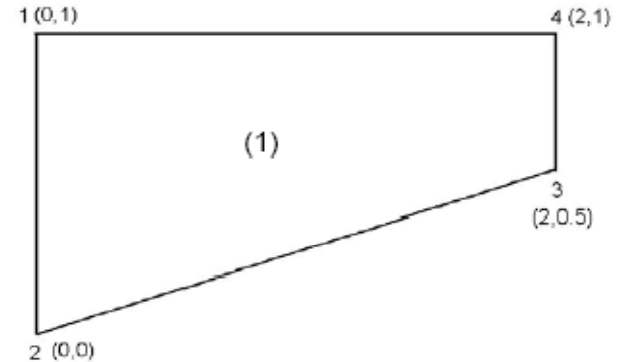


- The nodal force is (Gauss integration not shown):

$$\{f_{\Omega}^e\} = \int_{\Omega} f [N^e]^T d\Omega = \int_{-1}^1 \int_{-1}^1 f [N^e]^T |J^1| d\xi d\eta = \int_{-1}^1 \int_{-1}^1 6 \begin{bmatrix} \hat{N}_1^e \\ \hat{N}_2^e \\ \hat{N}_3^e \\ \hat{N}_4^e \end{bmatrix} (-0.125\eta + 0.375) d\xi d\eta = \begin{bmatrix} 2.5 \\ 2.5 \\ 2 \\ 2 \end{bmatrix}$$

Example problem using 4-node quadrilateral elements

- We now need to compute the contribution to element load from the natural boundary conditions on the element $e=1$ at $\xi=-1$.



$$\{f_q^e\} = - \int_{\Omega} \bar{q} [N^e]^T d\Gamma = - \int_{-1}^1 \bar{q} \left([N^e]^T |J^1| \right)_{\xi=-1} d\eta = - \int_{-1}^1 20 \begin{bmatrix} \frac{1}{2}(1-\eta) \\ 0 \\ 0 \\ \frac{1}{2}(1+\eta) \end{bmatrix} \underbrace{\frac{2-0}{2}}_{|J^1|} d\eta = \begin{bmatrix} -20 \\ 0 \\ 0 \\ -20 \end{bmatrix}$$

- Assembly of all the nodal contributions gives:

$$\{f\} = \{f_{\Omega}\} + \{f_{\Gamma}\} + \{r\} = \begin{bmatrix} r_1 - 17.5 \\ r_2 + 2.5 \\ r_3 + 2 \\ -18 \end{bmatrix}$$

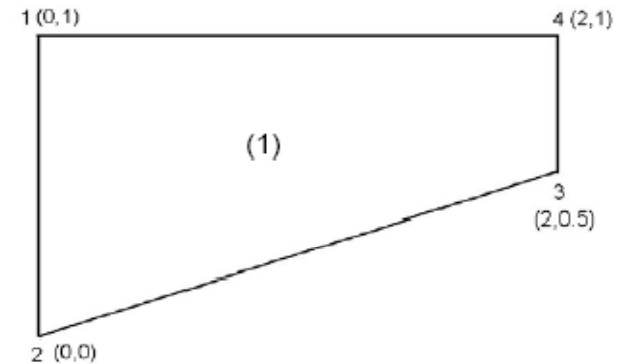
Example problem using 4-node quadrilateral elements

- We finally obtain:

$$\left[\begin{array}{ccc|c} 4.76 & -3.51 & -2.98 & 1.73 \\ -3.51 & 4.13 & 1.73 & -2.36 \\ -2.98 & 1.73 & 6.54 & -5.29 \\ \hline 1.73 & -2.36 & -5.29 & 5.91 \end{array} \right] \begin{bmatrix} 0 \\ 0 \\ 0 \\ T_4 \end{bmatrix} = \begin{bmatrix} r_1 - 17.5 \\ r_2 + 2.5 \\ r_3 + 2 \\ -18 \end{bmatrix}$$

which gives $T_4 = -3.04$.

- The nodal DOFs are: $\{d^1\} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ -3.04 \end{bmatrix}$



Example problem using 4-node quadrilateral elements

- The flux components at each Gauss point are:

$$\left\{ q^1(\xi, \eta) \right\} = \begin{bmatrix} -k \frac{\partial T^e}{\partial x} \\ -k \frac{\partial T^e}{\partial y} \end{bmatrix} = -k \left[B^e \right] \begin{bmatrix} T_1^e \\ T_2^e \\ T_3^e \\ T_4^e \end{bmatrix} = -k \begin{bmatrix} \frac{1+\xi}{3-\eta} & 1 \\ \frac{8}{\eta-3} & 0 \end{bmatrix} \frac{1}{4} \begin{bmatrix} \eta-1 & 1-\eta & 1+\eta & -\eta-1 \\ \xi-1 & -\xi-1 & 1+\xi & 1-\xi \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \\ -3.04 \end{bmatrix} \Rightarrow$$

$$\left\{ q^1\left(\frac{1}{\sqrt{3}}, -\frac{1}{\sqrt{3}}\right) \right\} = \begin{bmatrix} 0.8979 \\ 3.5916 \end{bmatrix}$$

$$\left\{ q^1\left(\frac{1}{\sqrt{3}}, \frac{1}{\sqrt{3}}\right) \right\} = \begin{bmatrix} -2.2965 \\ 19.793 \end{bmatrix}$$

$$\left\{ q^1\left(-\frac{1}{\sqrt{3}}, \frac{1}{\sqrt{3}}\right) \right\} = \begin{bmatrix} 4.9482 \\ 19.793 \end{bmatrix}$$

$$\left\{ q^1\left(-\frac{1}{\sqrt{3}}, -\frac{1}{\sqrt{3}}\right) \right\} = \begin{bmatrix} 5.8042 \\ 3.5916 \end{bmatrix}$$